

ENGG 225

David Ng

Winter 2017

Contents

1	January 9, 2017	5
1.1	Circuits, Currents, and Voltages	5
2	January 11, 2017	6
2.1	Ideal Basic Circuit Elements	6
3	January 13, 2017	8
3.1	Ideal Basic Circuit Elements Cont'd	8
3.2	Power and Energy	8
4	January 16, 2017	9
4.1	Power and Energy Cont'd	9
4.2	Kirchhoff's Laws	11
5	January 18, 2017	12
5.1	Kirchhoff's Laws Cont'd	12
6	January 20, 2017	13
6.1	Resistive Circuits	13
7	January 23, 2017	15
7.1	Circuit Analysis Using Series-Parallel Equivalentents	15
7.2	Other Simple Resistor Circuits - Voltage and Current Dividers	16
7.3	Node-Voltage Analysis	18
8	January 25, 2017	18
8.1	Node-Voltage Analysis Procedure	18
9	January 27, 2017	20
9.1	Node-Voltage Analysis Examples	20

10 January 30, 2017	22
10.1 Node-Voltage Analysis Examples Cont'd	22
10.2 Node-Voltage Method - Special Case	22
11 February 1, 2017	26
11.1 Mesh-Current Method	26
12 February 3, 2017	27
12.1 Mesh-Current Analysis Procedure	27
13 February 6, 2017	28
13.1 Mesh-Current Analysis Examples	28
13.2 Mesh-Current Analysis - Special Case	29
14 February 8, 2017	30
14.1 Supermesh Examples	30
14.2 Summary of Node-Voltage and Mesh-Current Methods	31
14.3 Thevenin and Norton Equivalent Circuits	31
15 February 10, 2017	31
15.1 Thevenin Circuits	31
16 February 13, 2017	32
16.1 Thevenin Circuit Examples	32
16.2 Shortcut Method for Thevenin Resistance	33
17 February 15, 2017	35
17.1 Thevenin Equivalent Circuits Conclusion	35
17.2 Principle of Superposition	36
18 February 27, 2017	37
18.1 Principle of Superposition Cont'd	37
18.2 Operational Amplifiers	38
18.3 Key Properties - Summing-Point Constraints	38
18.4 Applying Summing-Point Constraints	40
19 March 1, 2017	40
19.1 Op-Amp Circuits Cont'd	40
20 March 3, 2017	42
20.1 Op-Amp circuits Cont'd	42

21 March 6, 2017	43
21.1 Examples of Other Op-Amp Circuits	43
21.2 Input Resistance of Op-Amp Circuits	45
21.3 Another Application of Op-Amps - Comparators	46
22 March 8, 2017	46
22.1 Inductors and Capacitors	46
22.2 Power and Energy in the Capacitor	47
22.3 Capacitances in Series and Parallel	48
23 March 10, 2017	48
23.1 Capacitances in Series and Parallel Cont'd	48
23.2 Inductors	49
23.3 Power and Energy in the Inductor	50
23.4 Inductors in Series and Parallel	50
23.5 Steady State Sinusoidal Analysis	52
23.6 Sinusoidal Currents and Voltages	52
23.7 Root-Mean-Square Values	52
23.8 Relating to DC Circuits	53
24 March 15, 2017	54
24.1 Phasors	54
24.2 Complex Numbers Review	55
24.3 KVL and KCL Using Phasors	56
24.4 Summary of Phasor Summation Method	57
25 March 20, 2017	58
25.1 Phase Relationship Between Sinusoids	58
25.2 Complex Impedances	58
25.3 Summary of Impedances	60
25.4 Circuit Analysis With Phasors and Complex Impedances	60
26 March 22, 2017	61
26.1 Complex Impedance Examples	61
27 March 24, 2017	63
27.1 Thevenin Equivalent AC Circuits	63
28 March 17, 2017	64
28.1 Frequency Dependent Circuits	64
29 March 29, 2017	65
29.1 Superposition in AC Circuits	65
29.2 Power in AC Circuits	66

30 March 31, 2017	68
30.1 Power in AC Circuits Example	68
30.2 The Power Triangle and Other Power Relationships	69
31 April 3, 2017	71
31.1 DC Motors	71
31.2 Operating Characteristics of Motors	72
31.3 Speed Regulation	72
32 April 5, 2017	73
32.1 Electrical Circuit of DC Motors	73
32.2 Magnetization Curve	73
32.3 Power and Torque: Developed vs. Output	75
33 April 7, 2017	75
33.1 Shunt-Connected DC Machines	75
33.2 Separately Excited DC Machines	77
33.3 Permanent-Magnet DC Motors	77
33.4 Series-Connected DC Motors	77
33.5 Torque-Speed Characteristics	77
34 Monday April 10, 2017	78
34.1 Formulas for Final Exam	78

1 January 9, 2017

1.1 Circuits, Currents, and Voltages

The concept of electric charge is the basis for describing all electrical phenomenon. Charge exists in discrete quantities at integer multiples of 6.022×10^{-19} Coulombs. We note that this is the charge of one electron. An **electric circuit** is an interconnection of circuit elements connected in closed paths by conductors. The following are common components of circuits:

1. A **voltage source** is denoted by a circle encompassing a plus-minus sign. It is the supplier of energy.
2. A **resistor** is denoted by a zigzag line.
3. An **inductor** is denoted by a coil of wire. An inductor is used to store energy in the magnetic field.
4. A **capacitor** is denoted by a pair of plates. A capacitor is used to store energy in the electric field.
5. A **connection point** is denoted by a dot where the circuit elements meet.
6. A **conductor** is denoted by lines. These are most commonly wires.

Two fundamentally important electrical quantities are current and voltage. **Electric current** is the rate of flow of electric charge, and is given by

$$i(t) = \frac{dq(t)}{dt},$$

where $i(t)$ denotes the current in Amperes (A), $q(t)$ denotes the charge in Coulombs (C), and t denotes the time in seconds (s). That is, $1 A = 1 C/s$. Given $i(t)$, one can also find the total charge $q(t)$ by solving the integral

$$q(t) = \int_{t_0}^t i(t)dt + q(t_0).$$

We normally assign reference directions for current, each shown in a circuit diagram as an arrow in the indicated direction. We note that we can choose reference directions arbitrarily. For instance, if the current flow is actually in the opposite direction, then the value of i is simply opposite in sign.

We have **direct current** (DC) and **alternating current** (AC). In a graph of current versus time, a direct current is a constant value, whereas an alternating current takes a form similar to a sine wave. DC is used in energy sources such as batteries, while AC is used for house voltage. We make use of some common notation in our discussion of current. Around a circuit element A , the current flows

in direction i_A in an arbitrary direction. Around a circuit element A in between nodes a and b , the current flows in direction i_{ab} (from node a to node b), or in direction $i_{ba} = -i_{ab}$.

Voltage is the energy transferred to a circuit element per unit of charge flowing through it, and is given by

$$V(t) = \frac{dW(t)}{dq(t)},$$

where $V(t)$ denotes the voltage in Volts (V), $W(t)$ denotes energy in Joules (J), and $q(t)$ denotes charge. That is, $1 \text{ V} = 1 \text{ J/C}$. Voltages are assigned polarities to indicate the direction of energy flow. A diagram consisting of a “−”, followed by a circuit element A , followed by a “+” with i in the opposite direction indicates that energy is absorbed by A . When i flows in the same direction from left to right, then energy is supplied by A .

For analysis purposes, we assign arbitrary reference polarities to each circuit element, with the “−” and “+” in arbitrary positions on either side of the circuit element. If polarity is opposite, the value of V is then simply of opposite sign. We make use of some common notation in our discussion of voltage. Around a circuit element A , the voltage denoted by V_A with the positive and negative terminals on either side in an arbitrary order. Around a circuit element A in between nodes a and b , the voltage V_{ab} always has the first subscript positioned on the positive terminal. Thus, $V_{ba} = -V_{ab}$ has the positive terminal at node b as opposed to node a .

2 January 11, 2017

2.1 Ideal Basic Circuit Elements

Here, we will talk about conductors sources, and resistors. Later, we will bring in inductors and capacitors. All circuit elements are characterized by their **voltage-current** relationship.

Conductors are described by a blank rectangle with appropriately labelled current and voltage. It can also be expressed as a single line with the appropriately labelled current and voltage, with the additional note that $V = 0$. From this, we can define a conductor **short circuit**, which is between two points “shorted” together. The absence of a conductor between circuit elements is an **open circuit**

Sources can be categorized as independent voltage sources, dependent voltage sources, independent current sources, and dependent current sources.

Independent voltage sources are represented as independent voltage sources. Symbolically, they are represented with a current and a numerical value associated with voltage around a circle with + and −. A DC source has a constant voltage value (30 V for instance), whereas an AC source has a variable voltage ($100 \sin(120\pi t)$ for instance). We note the following properties for voltage sources:

- Voltage is specified explicitly. It is not dependent on any external factors.

- Voltage is unchanged by whatever it is connected to. For instance, voltage is independent of the current through it.

Dependent voltage sources have the same properties as independent sources, but the value of the voltage depends on either a voltage or a current elsewhere in the circuit. VCVS (voltage-controlled voltage source) are represented as a diamond shape with + and -. The voltage is written as μV_x where μ is the constant, and V_x is the controlling voltage. CCVS (current-controlled voltage source) on the other hand, are represented with the same shape, but with the current μi_x , where i_x is the controlling current. These can both be DC or AC as well. We note that in these cases, V_x and i_x are specified elsewhere in the circuit.

Independent current sources are represented with appropriately labelled voltage around a circle with an arrow pointing in one direction along a conductor. Its current could be labelled for instance, as 10 A , or $100 \cos(120\pi t)$. We note the following properties for current sources:

- Current is specified explicitly. That is, it is not dependent on external factors.
- Current is unchanged by whatever it is connected to. For instance, current is independent of the voltage across it.

Dependent current sources have the same properties as their independent counterparts, except their current depends on a voltage or current elsewhere. VCCS (voltage-controlled current source) is represented by a diamond shape with an arrow pointing towards the direction of the conductor. It is labelled as μV_x , where μ is a constant, and V_x is the controlling voltage. CCCS (current-controlled current source) is represented with the same shape, but with current μi_x , where i_x is the controlling current.

Resistors are represented with a zigzag within the conductor labelled R , with appropriately labelled current and voltage. Resistance is measured in Ohms (Ω). R is a constant. We note the property that voltage and current are related by **Ohm's Law**,

$$V = iR.$$

It is important to note that the direction of i and the polarity of V , are defined as shown for Ohm's Law. That is, the resistor is always absorbing energy. If we plot V on the y axis with i on the x axis, then the slope is equal to R . This results as a consequence of Ohm's Law. If we had instead labelled the resistor with current flowing from left to right, with voltage being - to +, then Ohm's Law states that $V = -iR$.

3 January 13, 2017

3.1 Ideal Basic Circuit Elements Cont'd

Conductance is related by Ohm's Law. Since $V = iR$, this can be rearranged so that

$$i = \left(\frac{1}{R}\right)V,$$

where conductance $G = \left(\frac{1}{R}\right)$. The SI units of conductance is Siemens (Ω^{-1}).

Remark. The unit of conductance was once mho (U)!

3.2 Power and Energy

Power is the product of voltage and current. That is,

$$P = Vi,$$

which we may also express as

$$P = \frac{dw}{dq} \times \frac{dq}{dt} = \frac{dW}{dt},$$

where P is the power in Watts (W), W is the energy in Joules, q is the charge in Coulombs, and t is the time in seconds. Thus, power is the rate of energy transfer.

We define power in terms of the **passive reference convention**. In this convention, current reference direction is the same direction as a voltage drop (from + to -). This implies that the circuit element absorbs power. For this scenario, $P = Vi$. If either current reference direction or voltage reference polarity is reversed, then we must use $P = -Vi$, which is the **active reference convention**. This is the case when current reference direction is in the direction of a voltage rise from - to +.

Example. Find the power in the circuit element given that the reference direction of i corresponds to the voltage rise from - to +. Given that $i = 10A$ and $V = 12V$, and then for $i = -10A$ and $V = 60V$.

We note that in either case, this corresponds to an active reference convention. Thus, we use the formula $P = -Vi$. In the first case, substituting values gives $P = -120W$. In the second case, we obtain $600W$.

Our physical interpretation of the sign of P is therefore that the circuit element absorbs power when $P > 0$, and the circuit element delivers power when $P < 0$. That is, in a circuit, depending on where + and - are placed with regards to each circuit element, since the current flows in one direction, we can use the formulas where $P = Vi$ and $P = -Vi$ to determine whether circuit element is delivering or absorbing energy.

In a resistor, we can apply our formulas to obtain an expression for power. By passive reference convention, we have $P = Vi$ and by Ohm's Law, we have $V = iR$. Substituting this expression of V , we obtain

$$P = (iR)i = i^2R.$$

We note that this is always positive, since a resistor is always absorbing power.

Energy can be determined by integrating the expression of power with respect to time. Since $P = \frac{dW}{dt}$, we get

$$W = \int_{t_1}^{t_2} P(t)dt + W(t_1).$$

Power companies measure energy to determine our monthly bills. The cost is determined by how much power is used over time.

Example. Suppose we are given a circuit with $i(t) = 2e^{-t}A$ flowing in the voltage direction from $+$ to $-$, where $V(t) = 10V$. Compute the power, compute energy from $t = 0 \rightarrow \infty$, and then determine whether energy is absorbed or delivered.

We recall that $P = Vi$. Therefore, $P(t) = (10V)(2e^{-t}A) = 20e^{-t}W$. To determine the energy consumed, we make use of the expression for energy. We obtain

$$\begin{aligned} W &= \int_0^{\infty} P(t)dt \\ &= \int_0^{\infty} 20e^{-t}dt \\ &= -20e^{-t} \Big|_0^{\infty} \\ &= 0 - (-20) \\ &= 20J \end{aligned}$$

Lastly, we note that W is positive, so the circuit element is absorbing energy.

4 January 16, 2017

4.1 Power and Energy Cont'd

In a circuit, depending on where $+$ and $-$ are placed with regards to each circuit element, since the current flows in one direction, we can use the formulas where $V = -iR$, $P = Vi$ and $P = -Vi$ to determine whether circuit element is delivering or absorbing energy. Suppose for instance that we know our power source to be at $100V$ and the resistor in the circuit to be at 10Ω . Then $V = -iR$, so $i = -\frac{V}{R}$. Substituting these values, we get $i = -\frac{100V}{10\Omega} = -10A$. In the source, we use passive reference convention to note that $P = Vi = (100V)(-10A) = -1000W$ delivered, while in the resistor, $P = -Vi = -(100V)(-10A) = 1000W$ consumed.

Example. Assume that energy cost is \$0.12 per kilowatt-hour (kWh). The electric bill for 30 days is \$60.00, with the power being constant over this time. Determine the power in watts. Given that voltage is 120V, determine the current. Lastly, determine how much energy is saved (in percent) by removing 60W.

First, we note that the energy consumed over the 30 days is

$$W = \$60.00 / \$0.12 = 500kWh.$$

Since power is constant over this time, it implies that on a power vs time graph, the use of power over the 30 days is constant. Thus, since power is the slope on an energy vs time graph, we have $W(t) = \int_0^t P dt = Pt$. Rearranging for power gives

$$\begin{aligned} P &= \frac{W}{t} \\ &= \frac{500kWh}{30 \text{ days}} \\ &= \frac{500000Wh}{30 \times 24h} \\ &= 694.4W \end{aligned}$$

Secondly, we want to determine the current through the circuit. Assuming that the house is absorbing energy, we have $P = Vi$, so

$$\begin{aligned} i &= \frac{P}{V} \\ &= \frac{694.4W}{120V} \\ &= 5.787A \end{aligned}$$

Lastly, if we reduce power consumption by 60W, this means we now save $60/694.4 * 100\% = 8.64\%$, where 8.64% of $\$60.00 = \5.18 .

Example. Consider the simple circuit with an independent voltage source of 15V and an independent current source of 2A. Furthermore, a current of 2A passes through the circuit from negative to positive on the voltage source and in the same direction as the current source. Determine the power in each source, and determine if the circuit element is absorbing or delivering power.

For the independent current element, we note that $P = Vi$, so $P_{2A} = (15V)(2A) = 30W$. For the independent voltage source, $P = -Vi$ since the current is traveling in the active reference convention, so $P_{15V} = -(15V)(2A) = -30W$. We note the energy balance that results since $-30W + 30W = 0$.

4.2 Kirchhoff's Laws

So far, we have reviewed fundamental electrical quantities of V , i , P , and W . We have also considered basic circuit elements, such as resistors and sources, each with their own unique $V-i$ relationship. Kirchhoff's laws can now be used to define how V and i are distributed in a circuit.

Kirchhoff's Current Law (KCL) states that the algebraic sum of all currents at a node must be zero. We choose a consistent way to distinguish between incoming and outgoing currents at a node. To understand Kirchhoff's Current Law, we can consider a fluid-flow analogy, whereby an incoming rate of 6 litres per minute combined with an incoming rate of 3 litres per minute results in an outgoing rate of 9 litres per minute. Analogously, we can consider a node in a circuit joining two or more circuit elements with incoming currents i_1 and i_2 , along with an outgoing current i_3 . Thus, incoming currents add and outgoing currents subtract. In our example, the sum of the currents is $i_1 + i_2 - i_3 = 0$. Therefore, $i_3 = i_1 + i_2$.

Remark. Note that two connection points joined by a conductor is equivalent to a single node. We can collapse this into a single connection point.

Example. Determine the outgoing current if there are incoming currents of $-5A$, $3A$, and $6A$.

According to KCL, at the node, we have the sum of the currents equal to zero. Therefore, $-5A + 6A + 3A - i_1 = 0$. Rearranging this, we find that $i_1 = 4$.

Series circuits are a very common and important circuit configuration. Let i_1 , i_2 , and i_3 be the currents through circuit elements 1, 2, and 3 respectively. Nodes A and B each join exactly two circuit elements, as they lie in between circuit elements 1 and 2, and circuit elements 2 and 3 respectively. According to KCL, we have at node A a current of $i_1 - i_2 = 0$, so $i_1 = i_2$. At node B , we have $i_2 - i_3 = 0$, so $i_2 = i_3$. Thus, we note that circuit elements in series must all have the same current. We cannot have circuits that violate KCL. We can distinguish circuit elements that are in series by examining all connection points (nodes), and identifying those where only two circuit elements are joined.

Kirchhoff's Voltage Law (KVL) states that the algebraic sum of all voltages around a loop must be zero. This law derives from conservation of energy. We consider circuit elements in a closed loop, where a loop is a closed path starting at a node and finishing back at the same node. We similarly need a consistent way to sum voltages. Around the loop, if we pass a circuit element from $+$ to $-$, then we add V . If we pass a circuit element from $-$ to $+$, then we subtract V . Circuits that violate KVL have voltages around a loop that do not sum to zero.

For instance, consider a closed loop where the current travels from $-$ to $+$ across circuit element A and B , and then from $+$ to $-$ across circuit element C . Applying KVL to this loop, we need $-V_A - V_B + V_C = 0$. This can be understood through relating KVL with power and energy, since there must be at all times an energy

balance (generated must equal absorbed). Hence at any time, the net power must be zero. Thus, we can consider circuit elements A , B , and C as having power $-V_A i$, $-V_B i$, and $V_C i$ respectively using passive reference convention. $-V_A i - V_B i + V_C i = 0$, or $i(-V_A - V_B + V_C) = 0$. Assuming that $i \neq 0$, this means that $-V_A - V_B + V_C = 0$, which is KVL.

5 January 18, 2017

5.1 Kirchhoff's Laws Cont'd

Parallel circuits are such that circuit elements in parallel have the same voltage. For instance, consider a circuit with circuit elements A , B , and C in parallel, with $+$ and $-$ for each in the same direction. We can then consider the loops that result when we compare these parallel elements. Around loop 1, we have $-V_A + V_B = 0$, so

$$V_A = V_B.$$

Around loop 2, we have $-V_B + V_C = 0$, so

$$V_B = V_C.$$

Therefore, we have

$$V_A = V_B = V_C$$

Example. Consider the circuit with circuit elements A , B , C , and D , where the first two circuit elements are in parallel with C and D (which are in series with each other). i_A flows upwards, i_B flows downwards, i_C flows upwards, and i_D flows downwards. V_A , V_B and V_C goes from $+$ to $-$ along with the direction of their current, whereas V_D goes from $+$ to $-$ against the direction of its current direction. Determine the circuit elements in series and the circuit elements in parallel. Determine i_C in terms of i_D . Lastly, given that $i_A = 3$ and $i_C = 1$, find i_B and i_D .

We note that only C and D are in series, since there is nothing also that joins where they join. That is, there is nothing else connected since the current cannot split at any point along the connection. Only A and B are in parallel, since there is nothing separating A and B , which are on separate paths each closed by two nodes. We note that they are not in parallel with either C or D , since they would either be separated by D or C in the respective cases. Since C and D are in series, they must have identical currents. Since i_D is in the opposite direction to i_C , we have $i_C = -i_D$. To determine i_B , we apply KCL where $i_A + i_C - i_B = 0$. Thus,

$$i_B = i_A + i_C = 3A + 1A = 4A.$$

As we have noted before, $i_C = -i_D$, so

$$i_D = -i_C = -1A.$$

Example. For the following circuit, let $V_0 = 100V$, and find the total power in the circuit using KVL and KCL. Let there be an unknown current i_g pointing downwards for this independent current source, followed by a voltage from $-$ to $+$ at $80V$. This is in parallel with an independent current source of $4A$ pointing upwards, with a current i_Δ and a voltage source from $+$ to $-$ of $80V$. We also have in parallel a voltage controlled voltage source V_0 from $-$ to $+$ with the current pointing upwards of $2i_\Delta$

We need to find all the voltages and currents. We first apply KCL at the location on top at node A , where the parallel paths meet. We note that

$$i_\Delta + 2i_\Delta - i_g = 0.$$

We know that $i_\Delta = 4A$, so we can isolate i_g and find that $4A + 2 \times 4A - i_g = 0$, so

$$i_g = 12A.$$

We can now redraw with our own labels for loops and unknown voltages. Let the leftmost loop be loop 1, and the rightmost be loop 2. The voltage across the independent current source of $12A$ is labelled V_{12} from $+$ to $-$ in the direction of the current, and the voltage across the independent current source of $4A$ is labelled V_4 from $-$ to $+$. Since we know all the currents, we now use this to find the unknown voltages V_{12} and V_4 . We first apply KVL around the rightmost loop to find that when we consider voltages across $+$ to $-$ as a positive voltage, with $-$ to $+$ indicate a negative voltage, we obtain $-V_4 + 80V + 100V = 0$, so

$$V_4 = 180V.$$

Now applying KVL to loop 1, we obtain $V_4 + 80V - V_{12} - 80V = 0$, so

$$V_{12} = 180V.$$

With the currents and voltages known, we can now determine the power using passive reference convention. For the $12A$ and $4A$ current sources, we have $2160W$ and $-720W$ respectively. For the independent voltage source from $-$ to $+$, we obtain $-960W$, whereas the voltage source from $+$ to $-$ has $320W$. Lastly, the dependent source has $-800W$. If we sum the total power in the circuit, we note that it equals $0W$. That is, there is an energy balance!

6 January 20, 2017

6.1 Resistive Circuits

KVL, KCL, and Ohm's Law give us all the tools we need to begin circuit analysis. We will now consider resistances in series and in parallel.

First we consider resistance in **series**. Consider an independent voltage source from $-$ to $+$ in the direction of current i . Along this series circuit is R_1 , R_2 , and R_3 resistors with voltages from $+$ to $-$ of V_1 , V_2 , and V_3 respectively. By KVL, we now have

$$-V + V_1 + V_2 + V_3 = 0.$$

We recall from Ohm's Law that for a current i across a resistor with voltage from $+$ to $-$, we have $V = iR$. Thus, the expression can be rewritten as

$$-V + iR_1 + iR_2 + iR_3 = 0.$$

In other words, we can isolate V and factor out i to get

$$V = i(R_1 + R_2 + R_3).$$

We note then that we can replace the resistors with a single equivalent resistance R_{eq} , where

$$R_{eq} = R_1 + R_2 + R_3.$$

Therefore, resistance in series add.

Now, we consider resistance in **parallel**. Suppose we have a parallel circuit with voltage source from $-$ to $+$ in the direction of current i . Along parallel paths, we have resistors R_1 , R_2 , and R_3 with currents i_1 , i_2 , and i_3 respectively with a common voltage V from $+$ to $-$. Then, we consider the common node and apply KCL to find that

$$i - i_1 - i_2 - i_3 = 0.$$

By applying Ohm's Law, we can rewrite this as

$$i - \frac{V}{R_1} - \frac{V}{R_2} - \frac{V}{R_3} = 0.$$

Therefore, solving for i by factoring out V , we obtain

$$\begin{aligned} i &= V \left(\frac{1}{R_1} + \frac{1}{R_2} + \frac{1}{R_3} \right) \\ &= V(C_1 + C_2 + C_3) \end{aligned}$$

Therefore, conductances in parallel add. In terms of voltage, this can be expressed as

$$V = i \left(\frac{1}{R_1} + \frac{1}{R_2} + \frac{1}{R_3} \right)^{-1}.$$

A very common resistor configuration consists of two resistors in parallel. The total resistance is therefore

$$\begin{aligned} R_{eq} &= \left(\frac{1}{R_1} + \frac{1}{R_2} \right)^{-1} \\ &= \frac{1}{\frac{1}{R_1} + \frac{1}{R_2}} \\ &= \frac{R_1 R_2}{R_1 + R_2} \end{aligned}$$

Example. Find a single equivalent resistance for a circuit with a 15Ω and 5Ω resistor in series. The circuit branches into two parallel paths, one with a 30Ω and 10Ω resistor in series, and the other path with a 40Ω resistor.

We note that we have two set of resistors in series. We can add these resistances, so we have 20Ω and 40Ω . We now evaluate the resistance of the two split paths by considering them in parallel. We apply the common expression of $R_{eq} = \frac{R_1 R_2}{R_1 + R_2}$ to obtain 20Ω . Since the result is now in series with the other 20Ω resistance, we simply add to obtain a final resistance of

$$R_{eq} = 20\Omega + 20\Omega = 40\Omega.$$

7 January 23, 2017

7.1 Circuit Analysis Using Series-Parallel Equivalents

Circuit Analysis is a procedure for determining all voltages and currents in every circuit element. We may employ the above simple resistor equivalents to analyze a circuit.

Example. Consider a circuit with an independent voltage source of $80V$ from $-$ to $+$, with a current i . This is met with a 60Ω resistor, followed by a 40Ω resistor on a parallel branch, and a 10Ω and 30Ω resistor on another parallel branch. Find the power in each of the circuit elements.

We first combine resistances. Since the 10Ω and 30Ω resistor are on the same branch, we can add the resistances to obtain 40Ω . Since this is in parallel with the other 40Ω resistor, we apply the expression for resistance in parallel to obtain 20Ω . This is now in series with the 60Ω resistor, so we add them to get a total resistance of $80V$. Since $V = iR$, we can isolate for current to obtain

$$\begin{aligned} i &= \frac{V}{R} \\ &= \frac{80V}{80\Omega} \\ &= 1A \end{aligned}$$

We can now reconstruct the original circuit and determine the voltage across each individual resistor. We first split back into the configuration with the 60Ω and 20Ω resistors (the 20Ω resistor is a combination of three resistors). Since $V = iR$, we have $60V$ and $20V$ respectively. We can check from KVL that this is correct since

$$-80V + 60V + 20V = 0.$$

We now split the 20Ω resistor into the remaining resistors. We note that the current splits off into two paths. Thus, we find the currents by once again applying the known voltage of $20V$ to each resistance. We find that for both paths, since the net resistance is 40Ω , $\frac{20V}{40\Omega} = 0.5A$. We can check KCL at the node where the current splits and note that this is correct since

$$1A - 0.5A - 0.5A = 0.$$

Now, we once again determine the voltage across each resistor. Over the 40Ω resistor with $0.5A$ we have the same $20V$, across the 10Ω resistor with $0.5A$ we have $5V$, and across the 30Ω resistor with $0.5A$ we have $15V$. We determine the power across each circuit element by considering all resistors use power, while the source generates the power. Thus, over the source we have $-80W$, over the first 60Ω resistor we have $i^2R = (1A)^2(60\Omega) = 60W$, over the 40Ω resistor we have $10W$, over the 10Ω resistor we have $2.5W$, and over the 30Ω resistor we have $7.5W$. Note the energy balance since

$$-80W + 60W + 10W + 2.5W + 7.5W = 0.$$

Later, we will use well established systemic methods to do analysis:

- **Node-Voltage Method**
- **Mesh-Current Method**
- **Thevenin Equivalentents**
- **Superposition**

7.2 Other Simple Resistor Circuits - Voltage and Current Dividers

In a series connection of resistors, the total applied voltage divides among them. We have $R_{eq} = R_1 + R_2 + R_3$, so

$$i = \frac{V}{R_{eq}} = \frac{V}{R_1 + R_2 + R_3}.$$

The individual voltages are therefore $V_1 = iR_1$, $V_2 = iR_2$, and $V_3 = iR_3$. For R_1 then, we have

$$V_1 = \left(\frac{R_1}{R_1 + R_2 + R_3} \right) V,$$

where the portion of the total resistance of R_1 is the same as the portion of V_1 of the total voltage. The results are similarly defined for R_2 and R_3 .

In a parallel connection of resistors, the total applied current divides among them. We have

$$R_{eq} = \left(\frac{1}{R_1} + \frac{1}{R_2} \right)^{-1} = \frac{R_1 R_2}{R_1 + R_2},$$

so

$$V = i \left(\frac{R_1 R_2}{R_1 + R_2} \right).$$

The individual currents are therefore

$$\begin{aligned} i_1 &= \frac{V}{R_1} \\ &= \frac{i}{R_1} \left(\frac{R_1 R_2}{R_1 + R_2} \right) \\ &= i \left(\frac{R_2}{R_1 + R_2} \right) \end{aligned}$$

where i_2 is defined analogously. Note the similarity to the voltage divider, except it is the resistor from the other branch in the numerator. It is not as straightforward as voltage division with more than two resistors. When more than two resistors are in parallel, we may group them.

Example. Find i_3 in the circuit with resistor 10Ω , which is in parallel with a 60Ω and 30Ω resistor with a $10A$ independent current source.

We recall that we can combine the other two resistors to obtain another resistor with equal net resistance. That is, $60\Omega * 30\Omega / (60\Omega + 30\Omega) = 20\Omega$. Now, applying the expression for current dividers, we obtain

$$\begin{aligned} i_3 &= i \left(\frac{R_2}{R_1 + R_2} \right) \\ &= 10A * \frac{20\Omega}{10\Omega + 20\Omega} \\ &= 6.67A \end{aligned}$$

Example. Given an independent current source of $2A$, with a 6Ω and 12Ω resistor in parallel, and a 12Ω and 24Ω resistor in parallel, find V , i_1 , and V_2 , where V is the total voltage generated by the current source, V_2 is the voltage across the 12Ω resistor in parallel with the 6Ω resistor, and i_1 is the current across the other 12Ω resistor.

There are many ways to proceed. For our purposes, we first find the total resistance connected to the source to find V . We then use voltage division to find V_2

and current division to find i_1 . We simplify the resistances to obtain the equivalent net resistance. By considering the parallel resistors, we obtain an 8Ω and 4Ω resistor. Combining these in series, we get a 12Ω resistor.

$$\begin{aligned}V &= iR \\ &= (2A)(12\Omega) \\ &= 24V\end{aligned}$$

Since circuit elements in parallel have the same voltage, we simply need to find V_2 by finding the voltage across the resistance over the net 4Ω resistor.

$$\begin{aligned}V_2 &= iR = \\ &= (2A)(4\Omega) \\ &= 8\Omega\end{aligned}$$

With the incoming total current of $2A$, we can apply current division to find i_3

$$\begin{aligned}i_3 &= i \left(\frac{R_2}{R_1 + R_2} \right) \\ &= 2A * \frac{24\Omega}{12\Omega + 24\Omega} \\ &= 1.33A\end{aligned}$$

7.3 Node-Voltage Analysis

The previous method of circuit analysis by series and parallel circuit manipulation works well for many circuits, but it is an “ad-hoc” method and depends on the circuit. Furthermore, it does not apply to all circuits. For instance, there are circuits where nothing is in series and nothing is in parallel. Node-voltage analysis works for any circuit. Its basic steps are as follows:

1. Identify nodes and decide on a reference node.
2. Apply KCL at nodes to develop a system of equations in terms of node voltages.
3. Solve for node voltages.

8 January 25, 2017

8.1 Node-Voltage Analysis Procedure

Consider a circuit with an independent voltage source V_s . The path breaks off at node A into two paths, one with resistor R_1 , and the other with resistor R_2 . The second path with R_2 continues to separate into two paths at node B , one with R_3

that joins with the path with R_1 at node C , and another with resistor R_4 . The joined path encounters R_5 before joining with the path with R_4 at node D , which then leads to the source.

We first identify the nodes, and decide on a reference node. We label node D as our reference node that we designate as our zero volt reference. All node voltages are relative to this reference node. It always simplifies the process by selecting a node at the side of a voltage source. We label this node with a short line extending from the node, followed by a long line, a short line, and a dot.

Note that according to this reference node, then node voltage V_1 at node A becomes $V_1 = V_s$. We now apply the second step and perform KCL at the nodes. So far, we have always labelled **branch voltages** for circuit elements. For instance, the current i_x across a resistor R_x from $+$ to $-$ has $V_x = i_x R_x$, and an independent voltage source has voltage V_x . However, we now need to consider **node voltages**, so we must express V_x and i_x in terms of node voltages V_1 and V_2 with V_2 the incoming voltage and V_1 the outgoing voltage.

According to common sense interpretation, voltage V_2 appears to be at a higher potential than V_1 . Therefore, the branch voltage V_x is the difference between the higher voltage V_2 and the lower voltage V_1 . That is,

$$V_x = V_2 - V_1.$$

In a KVL interpretation, we arrive at the same result by forming a loop with an independent voltage source and two resistors with voltage V_z , V_x , and V_y respectively. We know by KVL that

$$-V_y - V_x + V_z = 0.$$

However, we consider the nodes when the independent voltage source and the resistor meet as V_2 and the node where both resistors meet as V_1 . Now, we have $V_1 = V_y$ and $V_2 = V_z$, so

$$-V_1 - V_x + V_2 = 0,$$

and hence

$$V_x = V_2 - V_1.$$

Since we now have an expression for V_x , we can also derive an expression for current in terms of node voltages,

$$i_x = \frac{V_2 - V_1}{R_x}.$$

In our original circuit in consideration, we have V_1 at node A , V_2 at node B , and V_3 at node C . We can now consider V_2 in the original circuit, and write an expression for current over R_2 , R_3 , and R_4 ,

$$i_2 = \frac{V_2 - V_1}{R_2},$$

$$i_3 = \frac{V_2 - V_3}{R_3},$$

$$i_4 = \frac{V_2 - 0}{R_4}.$$

Remark. Note that by convention, we always point the arrow of the current away from the node of interest from the resistors from + to -. We have applied KVL and Ohm's Law to determine the current term in each case. Note that the expression for i_4 contains 0 since this is our reference node.

We now sum the currents at node B by KCL to obtain

$$-\left(\frac{V_2 - V_1}{R_2}\right) - \left(\frac{V_2 - 0}{R_4}\right) - \left(\frac{V_2 - V_3}{R_3}\right) = 0.$$

For simplicity, we multiply the entire equation by -1 to obtain

$$\left(\frac{V_2 - V_1}{R_2}\right) + \left(\frac{V_2 - 0}{R_4}\right) + \left(\frac{V_2 - V_3}{R_3}\right) = 0.$$

Repeating the above procedure for node C with currents once again pointing away from the node, we obtain

$$\left(\frac{V_3 - V_1}{R_1}\right) + \left(\frac{V_3 - V_2}{R_3}\right) + \left(\frac{V_3 - 0}{R_5}\right) = 0.$$

Remark. Note the general pattern when writing a node equation with a resistor branch. The node of interest comes first, followed by a subtraction of the connecting node, all over the connecting resistance. All terms in this equation represent the current leaving a node, summed to zero by KCL.

For objects other than resistors connected to the node, we perform different operations. For an independent current source leaving the node, we use the positive corresponding current which is a constant (independent current sources entering the node are therefore given the corresponding negative value). For an independent voltage source leaving the node from + to -, the current over this source is unknown. We employ a variation of the current method used.

9 January 27, 2017

9.1 Node-Voltage Analysis Examples

Example. An independent current source of $1A$ is connected to node A of voltage V_1 , splitting off into two paths. The first contains a 5Ω resistor, while the second contains a 2Ω resistor, which then branches off at node B with voltage V_2 into two paths. The first contains a 10Ω resistor that connects to the path with the 5Ω resistor at node C with voltage V_3 . The second path contains a 5Ω resistor that connects to node D , which is reached by the first branch after it passes a $10V$ independent voltage source from + to - and the reference node E . Determine V_1 , V_2 , and V_3 .

We note that since V_3 connects directly to the $10V$ source, $V_3 = 10V$. We can then consider node A with V_1 . By considering the outgoing currents, we obtain

$$\frac{V_1 - V_2}{2\Omega} + \frac{V_1 - 10V}{5\Omega} + (-1A) = 0.$$

We then consider node B with V_2 , where we obtain

$$\frac{V_2 - V_1}{2\Omega} + \frac{V_2 - 10V}{10\Omega} + \frac{V_2 - 0}{5\Omega}.$$

From both of these equations, we can solve the following equivalent system

$$7V_1 - 5V_2 = 30,$$

$$5V_1 + 8V_2 = 10.$$

Solving this system, we obtain $V_1 = 9.35V$ and $V_2 = 7.097V$.

Example. A 5Ω resistor connects to node A with voltage V_1 , splitting into two paths. The first contains a dependent current source of $2i_x$ pointing backwards. The second path connects to a 2Ω resistor, then to node B of voltage V_2 , with one path connecting to a 5Ω resistor with current i_x in the forward direction, connecting with the path of the dependent current source at node C of voltage V_3 . The second path connects to a $10V$ source which connects with the node. Solve for the node voltages.

We once again have 3 nodes, with $V_2 = 10V$ since it is connected to the voltage source. We therefore solve for V_1 and V_3 ,

$$\frac{V_1 - 0V}{5\Omega} + \frac{V_1 - V_2}{2\Omega} + (-2i_x) = 0,$$

$$\frac{V_3 - 0V}{10\Omega} + \frac{V_3 - V_2}{5\Omega} + 2i_x = 0.$$

Note that $\frac{V_3 - V_2}{5\Omega}$ is simply i_x pointing in the other direction. This expression is therefore equal to $-i_x$. Since we know $V_2 = 10V$, we can simplify the above equations to obtain

$$7V_1 - 50 - 20i_x = 0,$$

$$3V_3 - 20 + 20i_x = 0.$$

We further simplify this by considering that

$$i_x = \frac{V_2 - V_3}{5\Omega}.$$

Thus, we substitute this for i_x in the two equations to solve for V_1 and V_3 . Doing this, we obtain $V_1 = 1.43V$ and $V_3 = 20V$.

10 January 30, 2017

10.1 Node-Voltage Analysis Examples Cont'd

Example. Suppose we are given a 1A current source heading up to connect at V_1 , where the path splits into 3 paths at a node A with voltage V_1 . The path leading back down is connected to a 5Ω resistor, the path going across is connected to a 15Ω resistor, and the path above connects to a 10Ω resistor with a voltage of V_x from $-$ to $+$. The 15Ω resistor path splits into 2 at a node B with voltage V_2 , with one passing a $2V_x$ dependent voltage source from $+$ to $-$, which then connects to the reference node, and another path with a 10Ω resistor. This path connects with the path with other 10Ω resistor at a node C with a voltage of V_3 , and then splits off into a path with a 5Ω resistor and one with a 2A current source in the upwards direction.

In this problem, we have three voltages, where we know that $V_2 = 2V_x$ since the path from V_2 to reference node has a voltage of $2V_x$. At node A , we have

$$-1A + \frac{V_1 - 0V}{5} + \frac{V_1 - V_2}{15} + \frac{V_1 - V_3}{10} = 0.$$

We can rearrange this to sum the conductances connecting each node to obtain

$$V_1 \left(\frac{1}{5} + \frac{1}{15} + \frac{1}{10} \right) - V_2 \left(\frac{1}{15} \right) - V_3 \left(\frac{1}{10} \right) - 1 = 0.$$

At node C , we have

$$\frac{V_3 - V_1}{10} + \frac{V_3 - V_2}{10} + \frac{V_3 - 0V}{5} - 2 = 0.$$

Note that we also have the dependence at top, so we find that $V_x = V_3 - V_1$. Thus, since $V_2 = 2V_x$, this becomes $V_2 = 2(V_3 - V_1)$. We now solve for the unknown values of V_1 and V_3 in the system of equations to obtain

$$15V_2 - 7V_3 = 30,$$

$$V_1 + 2V_3 = 20.$$

We can then conclude that $V_1 = 5.405V$, $V_3 = 7.297V$, and $V_2 = 2(V_3 - V_1) = 3.784V$.

10.2 Node-Voltage Method - Special Case

There is only one special case that needs to be handled, where we have voltage sources between nodes where neither node is a reference node. First, the **simpler case** occurs when voltage sources are connected directly to other voltage sources. For instance, consider a reference node connected to 3 paths, each ending at nodes

with voltages of V_1 , V_2 and V_3 respectively. The first two paths have a resistor, while the third is connected to a voltage source of V_a from $-$ to $+$. Nodes 1 and 2 are connected by a resistor, and nodes 2 and 3 are connected by a voltage source V_b from $+$ to $-$. Nodes 1 and 3 are connected by another resistor. We notice that $V_3 = V_a$, and that $V_2 - V_3 = V_b$, so $V_2 = V_a + V_b$. In this case, V_2 and V_3 are already known, so we do not need to write equations. We are left with only one unknown with one equation to solve at node 1.

The **trickier case** occurs when the sources are not directly connected to each other. Suppose the same configuration as above, except V_b is now between V_2 and V_1 from $+$ to $-$ and V_2 and V_3 are connected by a resistor instead. While $V_3 = V_a$ is the same, we do not know the current i that spans the voltage source from V_1 to V_2 . Recall in writing node equations that we sum the currents leaving the nodes. At node 1, we have

$$\frac{V_1 - 0V}{R_1} + \frac{V_1 - V_3}{R_2} + i = 0,$$

where R_1 and R_2 are the resistances between the respective nodes. Since i is another unknown along with V_1 and V_2 , we cannot ignore it. We have to handle the problem with the concept of a **supernode** comprised of V_1 , V_2 , and the current voltage source in between. Considering node 2 of the original circuit, we have

$$\frac{V_2 - 0V}{R_3} + \frac{V_2 - V_3}{R_4} - i = 0,$$

where R_3 and R_4 are the resistances between the respective nodes. We can now eliminate i by adding the above two equations to obtain

$$\frac{V_1 - 0V}{R_1} + \frac{V_1 - V_3}{R_2} + \frac{V_2 - 0V}{R_3} + \frac{V_2 - V_3}{R_4} = 0.$$

This is the **supernode equation** where the left side of the supernode is the first two components of the expression, and the right side of the supernode is the remaining two components of the equation. We also have a **dependence equation** for the two nodes within the supernode

$$V_2 - V_1 = V_b.$$

The result is that we still have two equations to describe the nodes V_1 and V_2 : the supernode and dependence equations.

Example. *Suppose we have an independent voltage source of 10V connected to a reference node on one end, and to node V_1 on another. The path splits into one that connects to V_3 with 20Ω , and one that connects to V_2 with 10Ω . V_2 is connected to a path with 2Ω to the reference node and to V_3 with a voltage source of 5V from $-$ to $+$. V_3 is connected to the reference node with 5Ω resistance. Solve for node voltages.*

We form a supernode between V_2 and V_3 . This gives

$$\frac{V_2 - V_1}{10} + \frac{V_2 - 0V}{2} + \frac{V_3 - V_1}{20} + \frac{V_3 - 0V}{5} = 0.$$

From the diagram, we know that $V_1 = 10V$, so we can substitute to obtain the equation

$$12V_2 + 5V_3 = 30.$$

We also have the dependence equation of

$$V_3 - V_2 = 5V.$$

Simultaneously solving this system of equations gives $V_2 = 0.294V$ and $V_3 = 5.294V$.

Example. Suppose we have a circuit with a $2A$ current source pointing to node B . This leads to a 2Ω resistor to node A , which connects back to the current source, and to a 5Ω resistor that leads to node C . This leads to a path with a 8Ω resistor followed by a 2Ω resistor with voltage V_x from $+$ to $-$ which ends at node A , and two paths that lead to node D . The first path is across a $20V$ voltage source from $+$ to $-$, and the second path is across a 20Ω resistor. Node D leads to node A through an independent current source of $1A$ in the reverse direction. Determine V_x .

We can select different reference nodes. This is summarized below:

1. **Choice A:** We need equations for B , C , and D , where C and D form a supernode.
2. **Choice B:** We need equations for A , C , and D , where C and D form a supernode.
3. **Choice C:** We need equations for A and B , where D is fixed at $-20V$.
4. **Choice D:** We need equations for A and B , where C is fixed at $20V$.

The best choices appear to be C or D . However, we will choose to demonstrate this example with node A . Solving for node B , we obtain

$$-2A + \frac{V_B - 0V}{2\Omega} + \frac{V_B - V_C}{5\Omega} = 0.$$

The supernode equation and dependence equation are

$$\frac{V_C - V_C}{5} + \frac{V_C - 0V}{8 + 2} - 1 = 0,$$

$$V_C - V_D = 20V.$$

Solving this system of three equations with three unknowns, we find that $V_B = 4.71V$, $V_C = 6.47V$, and $V_D = -13.53V$. We now use a simple voltage divider to find that

$$V_x = \left(\frac{2}{8 + 2} \right) V_C = 1.294V.$$

Example. Suppose that we have a $0.25A$ to V_c . V_c is connected to reference through a resistor of 4Ω , to V_a with 1Ω , and to V_b with a dependent current source from of $4i_x$ from $-$ to $+$. V_b is connected to reference with 1Ω , and to V_a with a current source of $2A$ in the reverse direction. V_a is connected to reference through a 4Ω resistor with current i_a , and also through a 4Ω resistor and a $10V$ voltage source from $+$ to $-$. Determine the voltages V_a , V_b , and V_c .

At node a , we have

$$\frac{V_a - 10V}{4} + \frac{V_a - 0V}{4} + \frac{V_a - V_c}{1} + 2A.$$

At the supernode, we have

$$\frac{V_c - 0V}{4} - 0.25A + \frac{V_c - V_a}{1} + \frac{V_b - 0V}{1} - 2A,$$

which reduces to $-4V_a + 4V_b + 5V_c = 9$. In terms of supernode dependence, we know that $V_b - V_c = 4i_x$, where $i_x = \frac{V_a}{4}$. Substituting this into the dependence equation, we obtain

$$-V_a + V_b - V_c = 0.$$

We now have 3 equations to solve for the 3 unknowns. Solving this system of equations, we find that $V_c = 1V$, $V_a = 1V$, and $V_b = 2V$.

Example (A Wheatstone Bridge). Suppose we have a $15V$ source from $-$ to $+$ leading to node 1. Here, we split off into a path with a 1200Ω resistor which arrives at a node labelled a , and then connects to a 300Ω resistor which meets back with the other path. The other path is a 1000Ω resistor, connected to node b , which passes a resistance of R_4 before meeting the first path at node 2. This then connects back to the voltage source. A voltage V_{ab} spans from a to b with a labelled as $+$ and b labelled as $-$. Assume that this bridge is balanced so that $V_{ab} = 0V$. Determine R_4 . Now, set $R_4 = 200\Omega$ and connect a and b with a 250Ω resistor. Find the power in the 250Ω resistor.

This is essentially a pair of voltage dividers. When it is balanced, $V_{ab} = 0V$, so a and b have equal voltages. We find the voltage at a is

$$V_a = \frac{300}{200 + 1200} * 15V = 3V.$$

This can also be checked using the node-voltage method by letting node 2 be the reference node. Since the voltages $V_a = V_b$, we find that

$$V_b = \frac{R_4}{R_4 + 1000} * 15V = 3V.$$

Solving this gives $R_4 = 250\Omega$. For the second problem, we find that at node a , we have

$$\frac{V_a - V_c}{1200} + \frac{V_a - 0V}{300} + \frac{V_a - V_b}{250} = 0,$$

where $V_c = 15V$ is the voltage at node 1. At node b , we have

$$\frac{V_b - 15V}{1000} + \frac{V_b - 0}{200} + \frac{V_b - V_a}{250} = 0.$$

Solving this system, we find that $V_a = 2.817V$ and $V_b = 2.627V$. Thus, since $P = \frac{V^2}{R}$, we have

$$P_{250} = \frac{(V_a - V_b)^2}{250\Omega} = 0.144mW.$$

11 February 1, 2017

11.1 Mesh-Current Method

The **mesh-current method** is another useful systematic method of circuit analysis. To use this method, the circuit must be planar with no crossing conductors. **Mesh currents** can be imagined as currents circulating in a closed loop or mesh. We note that mesh currents are different from **branch currents** in that we use branch currents to write KCL equations, and mesh currents cannot be measured with an ammeter. For instance, consider the branch currents i_1 , i_2 , and i_3 , where i_1 is the only one that enters a node, and the other two leave the node. That is,

$$i_1 = i_2 + i_3.$$

In mesh currents with i_a and i_b around the parallel loops around the node, we have

$$i_a = i_1,$$

$$i_b = i_2,$$

$$i_3 = i_a - i_b.$$

The main steps for mesh-current analysis are as follows:

1. Identify meshes.
2. Write the mesh-current equation for each mesh to develop a system of equations.
3. Solve for mesh equations.

12 February 3, 2017

12.1 Mesh-Current Analysis Procedure

Consider a circuit with a voltage source V_a from $-$ to $+$ which leads to two paths at node A , one with a current source i_s in the forward direction leading to node C , and another with resistance R_1 to node B . At node B , a resistor R_3 connects to node D while a resistor R_2 connects to node C . Node C connects to node D through a voltage source of V_b from $+$ to $-$, and node D leads back to the negative end of the original voltage source.

To identify meshes, we can imagine a circuit as a window with panes. We then assign a mesh current to each pane. The bottom left mesh current between V_a , R_1 and R_3 is i_a , the mesh current between R_2 , R_3 , and V_b is i_b , and the mesh current between i_s , R_1 , and R_2 is i_c . We note immediately that i_s forces the mesh current $i_c = i_s$, so i_c is known immediately.

Secondly, we form mesh equations in each mesh. We do so by labelling the circuit to indicate a polarity on each resistor inside each mesh in response to the mesh current in that mesh. This means that in the direction of each mesh current, any resistor R_x goes from $+$ to $-$ in the direction of the mesh current being considered. R_1 for instance goes from $+$ to $-$ for both i_a and i_c when considering their respective cases, even when they are both approaching the resistor from opposite directions.

For mesh a , voltages around the mesh must sum to zero by KVL. To find the voltage across resistors R_1 and R_3 , we need to determine the branch currents in terms of mesh currents. For R_1 , the branch current i_1 is in the direction we chose for i_a , so

$$i_1 = i_a - i_c.$$

Since $V_1 = i_1 R_1$, we can substitute the branch current with the mesh currents to obtain

$$V_1 = (i_a - i_c)(R_1).$$

Summing the voltages around mesh a , we find that

$$-V_a + (i_a - i_c)R_1 + (i_a - i_b)R_3 = 0.$$

Similarly, for mesh b , we obtain

$$V_b + (i_b - i_a)R_3 + (i_b - i_c)R_2 = 0.$$

Now, we complete the third step by solving the system of two equations in the two unknowns of i_a and i_b since we already know $i_c = i_s$. We can now completely solve the circuit.

Example. Suppose we have a loop with a 1Ω resistor on the left, a 2Ω resistor on top, and a $5V$ source from $+$ to $-$ on the right. On top, the 2Ω resistor forms a loop with a 4Ω resistor and a 3Ω resistor. On the right, the $5V$ source forms a loop

with the 3Ω resistor and a $1A$ current source in the reverse direction. These loops have mesh currents of i_a , i_b , and i_c respectively. Find the power in the 3Ω resistor.

Note that we immediately know that $i_c = -1A$. We now have two unknowns of i_a and i_b . For mesh a , we have

$$(i_a)(1\Omega) + (i_a - i_b)(2\Omega) + 5V = 0.$$

Note that i_a is alone in the first term since there is only one mesh current in the 1Ω resistor. This is in contrast to the two mesh currents $i_a - i_b$ in the 2Ω resistor that are opposite in direction. Likewise, for mesh b we have

$$(i_b - i_a)(2\Omega) + (i_b)(4\Omega) + (i_b - i_c)(3\Omega) = 0.$$

Solving this system of equations with $i_c = -1A$ gives $i_a = -0.826A$ and $i_b = -2.217A$. To determine the power in the 3Ω resistor, we note that the branch current in the direction of i_c is $i_c - i_b = -0.174A$. Thus, we find power to be

$$P = i^2 R = (-0.174A)^2 (3\Omega) = 0.0908W.$$

13 February 6, 2017

13.1 Mesh-Current Analysis Examples

Example. Suppose we have a circuit with a dependent voltage source of $10i_x$ from $-$ to $+$ followed by a 10Ω resistor and a 5Ω resistor in mesh current i_a , where i_x is the branch current on the branch with 5Ω in the direction of i_a on that branch. i_b consists of the 5Ω resistor with a 20Ω resistor and a $10V$ source from $+$ to $-$. Determine the mesh currents i_a and i_b .

For mesh a and b , we have

$$-10i_x + 10i_a + 5(i_a - i_b) = 0,$$

$$5(i_b - i_a) + 20i_b + 10 = 0.$$

We can express i_x in terms of the mesh currents, so

$$i_x = i_a - i_b.$$

Thus, solving this system of equations, we obtain $i_a = \frac{1}{3}A$, and $i_b = -\frac{1}{3}A$.

Example (Wheatstone Bridge). Suppose we are given a circuit with a $15V$ independent voltage source from $-$ to $+$ in the direction of mesh current i_a . This leads to node A , where we find resistances of 1000Ω , 250Ω , and 1200Ω forming i_b between nodes A , B , and C . This is followed by mesh i_c comprised of resistances of 250Ω , 200Ω , and 300Ω between nodes C , B , and D , with node D connecting back to the voltage source. Note that i_a consists of the voltage source, the 1200Ω resistor, and the 300Ω resistor. Determine the power in the 250Ω resistor.

We list the mesh current equations for i_a , i_b , and i_c respectively,

$$-15 + 1200(i_a - i_b) + 300(i_a - i_c) = 0,$$

$$250(i_b - i_c) + 1200(i_b - i_a) + 1000i_b = 0,$$

$$300(i_c - i_a) + (250(i_c - i_b) + 200i_c) = 0.$$

Note that similar to the node-voltage method, we can see an important pattern arise when we rearrange the equations. Consider the rearrangement of the equation for mesh a :

$$(1200 + 300)i_a - 1200i_b - 300i_c - 15 = 0.$$

These terms represent the total resistance around mesh a , the resistance shared with mesh b , the total resistance shared with mesh c , and the total voltage respectively. The three equations above form a system with three unknowns. The branch current in the 250Ω resistor in the direction of i_c is given by $i_c - i_b$. Thus, since $P = i^2R$, this becomes

$$P = (i_c - i_b)^2(250).$$

13.2 Mesh-Current Analysis - Special Case

As in the node-voltage method, there is a special case that we must handle. This occurs when there is a branch current source shared between meshes. Suppose we are given a circuit where R_1 , R_2 , and the independent current source i_s in the forward direction form i_a , the independent current source i_s in the opposite direction, R_4 , and an independent voltage source V_b from $+$ to $-$ forms i_b , and R_2 , R_3 and R_4 form i_c . We note that i_s is shared by meshes a and b .

In mesh a , we have an unknown voltage V because of the current source, so we obtain

$$i_a R_1 + (i_a - i_c)R_2 + V = 0.$$

In mesh b , we also have the unknown voltage V ,

$$-V + (i_b - i_c)R_4 + V_b = 0.$$

Now, we can add these two equations together to obtain the **supermesh equation** by considering a and b a supermesh,

$$i_a R_1 + (i_a - i_c)R_2 + (i_b - i_c)R_4 + V_b = 0.$$

The first two terms represent the side of the supermesh in mesh a , while the last two terms represent the side of the supermesh in mesh b . We also have a **dependence equation** for the two meshes within the supermesh,

$$i_a - i_b = i_s.$$

Example. Suppose we have a circuit with a $10V$ voltage source from $-$ to $+$ with a 5Ω resistor and a $2V_x$ current source in the opposite direction forming i_a . We also have this $2V_x$ current source in the forward direction, with a 10Ω resistor with a voltage of V_x from $+$ to $-$ and a $5V$ voltage source from $+$ to $-$ forming i_b . Determine the mesh currents.

We need a supermesh equation and a dependence equation. These are given respectively as

$$\begin{aligned} -10 + 5i_a + 10i_b + 5 &= 0, \\ i_b - i_a &= 2V_x. \end{aligned}$$

Taking into account the dependent current source, we note that $V_x = 10i_b$. Solving this system gives $i_a = 1.118A$ and $i_b = -0.0588A$.

14 February 8, 2017

14.1 Supermesh Examples

Example. Suppose we have mesh i_1 with a 10Ω resistor connected to a $3A$ current source in the forward direction, a 5Ω resistor and a $2A$ current source in the reverse direction. Mesh i_2 consists of the same $2A$ current source in the forward direction, a 15Ω resistor, and a 10Ω resistor. Mesh i_3 consists of the same 15Ω resistor and 5Ω resistor, along with another 5Ω resistor and a 10Ω voltage source from $-$ to $+$. Determine the mesh currents and find the power in the current sources.

We note that we have a supermesh in which $i_1 = 3A$ is already known. From the supermesh dependence equation, we have

$$i_2 - i_1 = 2A,$$

so $i_2 = 5A$. For i_3 , we have

$$5i_3 - 10V + 15(i_3 - i_2) + 5(i_3 - i_1) = 0.$$

Solving for i_3 , we obtain $i_3 = 4A$. Now to determine power, we need to find the unknown voltages V_{2A} and V_{3A} from $-$ to $+$ in the forward direction across the $2A$ and $3A$ current sources respectively.

Remark. ASK HOW THE DIRECTION OF VOLTAGE IS ASSIGNED.

For i_2 , we have

$$-V_{2A} + 15(i_2 - i_3) + 10i_2 = 0,$$

so $V_{2A} = 65V$. The power in the $2A$ current source is therefore $P = -(65V)(2A) = -130W$. For i_1 , we have

$$-V_{3A} + 5(i_1 - i_3) + V_{2A} + 10(i_1) = 0,$$

so $V_{3A} = 90V$. Likewise, the power over this current source is $P = -(90V)(3A) = -270W$. In both cases, we note that the power is supplied.

14.2 Summary of Node-Voltage and Mesh-Current Methods

When choosing between Node-Voltage and Mesh-Current methods, we pick the method with fewer equations. For Node-Voltage, this means looking for nodes with voltage sources attached. This may eliminate equations through a good choice of reference node. For Mesh-Current, this means looking for meshes where mesh currents are fixed in value by current sources. So far, we have covered Circuit Simplification, KVL, KCL, Ohm's Law, Node-Voltage, and Mesh-Current. We now proceed to cover Thevenin Theorem and the Principle of Superposition.

14.3 Thevenin and Norton Equivalent Circuits

Theorem (Thevenin's Theorem). *A DC electrical network containing voltage sources, current sources, resistors, and two terminals is electrically equivalent to a network with one voltage source and one resistor.*

This gives us a way to arbitrarily complex "two-terminal" circuits, modeled as Thevenin voltage and Thevenin resistance, denoted as V_t and R_t respectively.

15 February 10, 2017

15.1 Thevenin Circuits

The voltage-current characteristics are identical at the two terminals. V_t and R_t are found by considering two operating extremes. In an open circuit comprised of independent and dependent sources and resistors, the current leading from the negative to the positive node is $i = 0$, where V_{DC} is the voltage across. In the Thevenin equivalent, we can replace the circuit with an independent current source V_t with a resistor R_t , with a current of $i = 0$ leading from the negative to the positive nodes. The voltage across the positive and negative node is V_t , so

$$V_t = V_{DC}.$$

In a short circuit comprised of independent and dependent sources and resistors, the current is i_{sc} , which forms a closed loop. The Thevenin equivalent would replace the sources with an independent voltage source V_t , and the resistors with R_t . Thus, since $V = iR$, we can obtain the following expressions for current and resistance,

$$i_{sc} = \frac{V_t}{R_t} = \frac{V_{DC}}{R_t},$$

$$R_t = \frac{V_{DC}}{i_{sc}}.$$

Determining a Thevenin equivalent is two separate analysis problems. We first need to find V_t , and then find R_t .

Example. Let there be a node x , which connects to node b , which leads to a 20Ω resistor before reaching node y . Node b also connects to a 15Ω resistor that leads to node a . Node a leads to a 10Ω resistor that leads to y , and also leads to a path with a 5Ω resistor connected to a $10V$ voltage source from $+$ to $-$ that passes to a reference node before reaching node y . Find the Thevenin equivalent at terminals x , y .

We first find V_t , where $V_t = V_{DC}$. Using the node-voltage method, the open-circuit voltage will be V_b . Thus, at node a ,

$$\frac{V_a - 10}{5} + \frac{V_a}{10} + \frac{V_a - V_b}{15} = 0.$$

At node b , we have

$$\frac{V_b - V_a}{15} + \frac{V_b}{20} = 0.$$

Solving this system of equations, we find that $V_a = 6.087V$ and $V_b = 3.478V$. Thus,

$$V_t = V_b = V_{DC} = 3.478V.$$

Now, we must find i_{sc} . Since there is a parallel combination with the branch containing the 20Ω resistor, we note that the equivalent resistance with the short circuit is 0Ω . V_b is now connected directly to the reference node, so $V_b = 0$. Similarly, we note that i_1 , the current through the 20Ω resistor is 0, since $V_b/20 = 0$. Thus, all current is in i_{sc} , with none in the 20Ω resistor. We can thus redraw the circuit. We have eliminated node b , and are left with a current i_{sc} on the branch containing the 15Ω resistor. The remaining parts of the circuit remain unchanged. We find again the voltage at node A ,

$$\frac{V_a - 10}{5} + \frac{V_a}{10} + \frac{V_a}{15} = 0.$$

Note above that $V_a/15$ is i_{sc} . Solving this, we find that $V_a = 5.455V$. Thus, ,

$$i_{sc} = \frac{V_a}{15} = 0.364A,$$

$$R_t = \frac{V_{DC}}{i_{sc}} = \frac{3.478V}{0.364A} = 9.56\Omega.$$

Therefore, the Thevenin equivalent circuit is node y , followed by $V_t = 3.478$ from $-$ to $+$, and a resistance $R_t = 9.56\Omega$ that leads to node x .

16 February 13, 2017

16.1 Thevenin Circuit Examples

Example. Let there be a node x which leads to node b . There is a $1A$ current source in the reverse direction to node y , and a 5Ω resistor to node y from node b . Node

b also leads to a 20Ω resistor before reaching node a . Node a is connected to node y through a $2A$ current in the reverse direction, and a 10Ω resistor. Node a is also connected to a 10Ω resistor that leads to a voltage source from $+$ to $-$, which crosses a ground before reaching y . Find the Thevenin equivalent.

We will find V_t using the node-voltage method, where $V_t = V_b$. At node a , we have

$$\frac{V_a - 10}{10} - 2 + \frac{V_a}{10} + \frac{V_a - V_b}{20} = 0.$$

At node b , we have

$$\frac{V_b - V_a}{20} + \frac{V_b}{5} - 1 = 0.$$

Solving this system for V_b gives $V_t = V_b = 6.667V$. We now need to find i_{sc} . Generally, when we solve for i_{sc} by connecting the nodes x and y , we usually treat this circuit as a new analysis problem. In this case, node b is again attached to the reference node, so $V_b = 0$. We once again perform node-voltage on nodes a and b to find,

$$\begin{aligned} \frac{V_a - 10}{10} - 2 + \frac{V_a}{10} + \frac{V_a - 0}{20} &= 0, \\ \frac{0 - V_a}{20} + \frac{0}{5} - 1 + i_{sc} &= 0. \end{aligned}$$

Note above that the 5Ω resistor is shorted out since there is no current there. We find that $i_{sc} = 1.6A$ by solving the system. Therefore,

$$R_t = \frac{V_t}{i_{sc}} = \frac{6.667V}{1.6A} = 4.1667\Omega.$$

The Thevenin equivalent is node x connected to a resistor of $R_t = 4.1667\Omega$, followed by a voltage source from $+$ to $-$ of $V_t = 6.667V$ that ends at node y .

16.2 Shortcut Method for Thevenin Resistance

If a circuit has no dependent sources, then we may use an alternative method to find R_t by zeroing the sources. We zero their values and use their “effective” resistance. For instance, suppose we have nodes a and b connected by an independent voltage source from $+$ to $-$. If we let $V = 0$, the source becomes a short circuit. That is, the effective resistance becomes 0Ω . If we have an independent current source instead from b to a , we can let $i = 0$, so the source becomes an open circuit, The effective resistance would therefore become $\infty\Omega$.

Example. Consider the example directly above. Using $V_t = 6.667V$, zero the sources to find R_t .

We note that we have removed the voltage source and replaced it with a short circuit, and replaced the two current sources with open circuits. This becomes a series-parallel combination of resistors. We find that the resistance is $R_t = 4.1667\Omega$ when we resolve this.

The above short-cut method cannot be used if a circuit has dependent sources. When there are dependent sources, we must determine i_{sc} to evaluate $R_t = V_t/i_{sc}$.

Example. Let node x be connected to node a through a 20Ω resistor. This leads to y through a path with a 15Ω resistor with current i_x , and alternatively on a path with a $2A$ independent current source in the reverse direction. Node a is also connected through a 10Ω resistor and a $5i_x$ voltage source from $+$ to $-$ to ground, which leads to y . Find the Thevenin equivalent.

First, we find V_t . With the terminals x and y in open circuit, no current flows through the 20Ω resistor, so $V_t = V_{DC} = V_a$. At node a , node-voltage states that

$$\frac{V_a - 5i_x}{10} - 2 + \frac{V_a}{15} = 0.$$

We determine V_a by solving the system formed by this equation with the knowledge that $i_x = V_a/15$ from the location of i_x . We find that $V_a = V_t = V_{DC} = 15V$. Now, i_{sc} can be found by considering the closed circuit. In this case, $i_{sc} = V_a/20$. By applying node-voltage, we find that

$$\frac{V_a - 5i_x}{10} - 2 + \frac{V_a}{15} + \frac{V_a}{20} = 0.$$

Once again equating this with $i_x = V_a/15$, we find that $V_a = 10.91V$ in this second scenario. Therefore,

$$i_{sc} = \frac{V_a}{20} = 0.545A,$$

$$R_t = \frac{V_t}{i_{sc}} = \frac{15V}{0.545A} = 27.5\Omega.$$

The Thevenin equivalent therefore consists of nodes x and y connected by a 27.5Ω resistor and a $15V$ voltage source.

Example. Let node x be connected to node c . Node c is connected to node b through a 15Ω resistor, to a through a 10Ω resistor with voltage V_x from $+$ to $-$, and to ground leading to y through a 20Ω resistor. Node b is connected to ground to y through a $2V_x$ current source in the reverse direction, and to node a through a 5Ω resistor. Node a is connected to a $10V$ voltage source from $+$ to $-$ leading through ground to y . Find the Thevenin equivalent.

We need to first solve for $V_t = V_{DC} = V_c$. Note that $V_a = 10V$. Solving for node-voltage at b and c , we find respectively,

$$\frac{V_b - 10}{5} - 2V_x + \frac{V_b - V_c}{15} = 0,$$

$$\frac{V_c - V_b}{15} + \frac{V_c}{20} + \frac{V_c - 10}{10} = 0.$$

For the dependent current source, we have

$$V_x = V_a - V - c = 10 - V_c.$$

We find that $V_c = V_t = V_{DC} = 9.29V$. We now find R_t . Since there is a dependent source, there are no shortcuts. Thus, the 20Ω resistor is short circuited. There is no voltage across, so there is also no current through it. Node-voltage through b is,

$$\frac{V_b - 10}{5} - 2(V_a - V_c) + \frac{V_b - V_c}{15} = 0.$$

However, we have established that $V_a = 10V$ and now $V_c = 0V$. We find that $V_b = 82.5V$ from this equation. Applying node-voltage to node c ,

$$\frac{V_c - V_b}{15} + \frac{V_c - 10}{10} + i_{sc} = 0.$$

Substituting $V_b = 82.5V$ and $V_c = 0V$, we find that $i_{sc} = 6.5A$. Thus,

$$R_t = \frac{V_t}{i_{sc}} = \frac{9.29V}{6.5A} = 1.43\Omega.$$

17 February 15, 2017

17.1 Thevenin Equivalent Circuits Conclusion

The **Norton equivalent** circuit provides an alternative form to the Thevenin equivalent. Instead of specifying R_t and V_t in series where

$$V_{DC} = V_t,$$

$$i_{sc} = \frac{V_t}{R_t},$$

the Norton equivalent replaces this with a current source i_n in parallel with a resistance R_t , where

$$V_{DC} = i_n R_t = V_t,$$

$$i_{sc} = i_n = \frac{V_t}{R_t}.$$

Thevenin and Norton equivalents are related by a source transformation (from voltage source and series resistance to current source and parallel resistance and vice versa).

Example. Consider the Thevenin equivalent composed of the node x connected to a resistance of $R_t = 1.43\Omega$, and a voltage source of $V_t = 9.29V$ leading to node y . Determine the Norton equivalent.

The Norton equivalent would have a parallel resistance of $R_t = 1.43\Omega$ with a current source of $i_n = V_t/R_t = 6.5A$.

Source transformations act as a handy simplification. Circuits can often be simplified by source transformations. Whenever we only have a voltage source from $-$ to $+$ followed by a resistance in series, we can replace it with a current source equalling $i = V/R$ along with the original resistance in parallel.

Example. *Simplify the circuit consisting of node a to ground. There are three paths from a to ground. The first path is across a 10Ω resistor followed by a $10V$ voltage source. The second path is across a 20Ω resistor followed by a $20V$ voltage source. The third path is across a 20Ω resistor followed by a $10V$ voltage source.*

We can replace the first two voltage source with a $1A$ current source, and the last voltage source with a $i = V/R = 10V/20\Omega = 0.5A$ current source. All circuit elements are now in parallel. We can now determine the total current, which can be replaced with a $2.5A$ current source. The resistors in parallel give an effective resistance of 5Ω . Thus, the final circuit becomes a $2.5A$ current source connected to a 5Ω resistor. The voltage across the resistor is

$$V = iR = 2.5A(5\Omega) = 12.5V.$$

17.2 Principle of Superposition

This is a fundamentally important concept, and often a required method in AC circuit analysis. We first present this for DC circuits. The method is as follows:

1. Let only one independent source be active.
2. Zero all other independent sources.
3. Determine the response r' (voltage or current) at the desired location in the circuit.
4. Repeat one at a time for all other independent sources in the circuit. That is, find r'' , r''' , etc.

The **principle of superposition** states that the total response r (voltage or current), is the sum of the individual responses,

$$r = r' + r'' + \dots + r^{(n)}.$$

Remark. Recall that removing current sources leaves an open circuit, while removing a voltage source leaves a short circuit.

Example. *Let there be a $50V$ voltage source from $-$ to $+$ with current i_x . This splits into two paths at node a with one leading through a 20Ω resistor to node c ,*

and another leading through a 12Ω resistor to node b . Node b leads to node c through an 8Ω resistor that connects to a $30V$ voltage source from $-$ to $+$ before reaching the $50V$ source. Node b also leads to the $50V$ source through a $20A$ current source in the reverse direction. Find i_x by superposition.

First, we will consider the $50V$ source by itself. The 12Ω and 8Ω resistors are now in series. This is in parallel with a 20Ω resistor, so the effective resistance is 10Ω . Thus,

$$i'_x = \frac{V}{R} = \frac{50V}{10\Omega} = 5A.$$

Now, considering the $20A$ source by itself, we note that the endpoints for the 20Ω resistor at nodes a and c are connected together. This means that $V_a = V_c$, so there are $0V$ and $0A$ across this resistor. Thus, we ignore this 20Ω resistor. We therefore have a circuit with a $20A$ current source with resistances of 12Ω and 8Ω . We use a current divider to find $-i_1$ in the loop with the 12Ω resistor, in order to find i''_x , since

$$i''_x = i_1 = -\frac{8}{8+12} * 20A = -8A.$$

Lastly, considering the $30V$ source by itself, we simply have the $30V$ source along with three resistors. The total resistance is 10Ω , so the current is

$$i'''_x = \frac{V}{R} = \frac{30V}{10\Omega} = 3A.$$

Thus, by the principle of superposition,

$$\begin{aligned} i_x &= i'_x + i''_x + i'''_x \\ &= 5A - 8A + 3A \\ &= 0A \end{aligned}$$

18 February 27, 2017

18.1 Principle of Superposition Cont'd

With dependent sources, we can only zero independent sources.

Example. SA $10V$ source from $-$ to $+$ is connected to a 5Ω resistor through a current i_x . This leads to a node that branches to a 10Ω resistor and a $2i_x$ voltage source from $+$ to $-$, and also branches to a $2A$ current source in the reverse direction with voltage V_x from $+$ to $-$. Find V_x .

Considering the $2A$ source alone, we use node voltage at node a where they connect, and designate the ground as the other end of the $2A$ source. Thus,

$$\frac{V_a}{5} - 2 + \frac{V_a - 2i'_x}{10} = 0.$$

Additionally, we know that

$$i'_x = -\frac{V_a}{5},$$

so solving this gives $V_a = 5.88V$. Thus, $V'_x = 5.88V$. Now, we consider the 10V source acting alone. Once again considering node a , we have

$$\frac{V_a - 10}{5} + \frac{V_a - 2i''_x}{10} = 0.$$

We also note that

$$i''_x = \frac{10 - V_a}{5}.$$

Solving this system gives $V_a = V''_x = 7.06V$. Finally, by superposition, we have

$$\begin{aligned} V_x &= V'_x + V''_x \\ &= 5.88V + 7.06V \\ &= 12.94V \end{aligned}$$

18.2 Operational Amplifiers

An **operational amplifier (op amp)** is a complex electronic circuit that implements a voltage-controlled voltage source. There are many important practical engineering examples including high-speed video amplifiers, microelectronic filters (telecommunications - huge industry), and instrumentation (precision measuring devices). Invented in 1968, it was originally used to perform operations in “analog” computers. It was used to perform operations such as addition, integration, and multiplication on voltages and currents.

Starting at node a , the non-inverting terminal, we reach a resistance R_i , and then reach node b , the inverting terminal. Node a is $+$, while node b is $-$, with a voltage of V_d across. Now, we have another starting at node c , passing through a resistor of $R_0 = 50\Omega$ (not 0Ω), a dependent voltage source of AV_d from $+$ to $-$, leading to node d . $R_i \approx 1M\Omega$ (not ∞), and $A = 100000$ (not ∞).

Since the op-amp is an electronic circuit, it requires an external source to operate. An example would be a ground connection from both end nodes, each leading to a voltage source from $-$ to $+$, with one holding a positive voltage, and the other a negative voltage. The path from the one holding the negative voltage meets the other path through a triangle from $-$ to $+$. the triangle has a vertex extending to a node, and two nodes extending from the side opposite to this vertex.

18.3 Key Properties - Summing-Point Constraints

1. **Virtual Short Circuit:** For any practical circuit, V_0 must be finite-valued, so $\|V_0\| < \infty$. For an ideal op-amp, $A = \infty$, and

$$V_0 = AV_d.$$

Thus, we can rearrange this to find that $V_0/A = V_d$. Since $A = \infty$, this gives

$$V_d = 0.$$

That is, there is no voltage across input terminals.

2. **Virtual Open Circuit:** For an ideal op-amp, $R_i = \infty$.

The circuit symbol for op amps is shown as a two ground connections connected to independent voltage sources V_1 and V_2 , both from $-$ to $+$, leading to the op amp. The path leading to the negative node on the side of the triangle is the **inverting input (-)**. The path leading to the positive node on the side of the triangle is the **non-inverting input (+)**. The vertex of the triangle leads to a positive node. This is not connected to another ground that exists leading from a negative node. The voltage across this positive and negative node is $V_0 = A(V_1 - V_2)$.

the op amp amplifies the differential input voltage V_d , where

$$V_d = V_1 - V_2 \rightarrow V_0 = AV_d,$$

where A is a large number. In the model of the ideal op amp, we have currents i_p and i_n coming in from the paths containing V_1 and V_2 respectively. Ideally, inputs look like open circuits ($i_n = 0, i_p = 0$). Inside the op amp is a voltage controlled voltage source where $V_0 = AV_d$, connected from ground, leading to $-$ to $+$, which leads out the vertex of the triangle.

The main characteristics of an ideal op amp are:

- Infinite input resistance, so $i_n = i_p = 0$. That is, no current flows into input terminals.
- $A = \infty$ (A is called the **open-loop gain**).
- Zero output resistance (the effective resistance of a voltage source is 0Ω).

In summary, the **summing-point constraints** are:

$$V_d = 0,$$

$$i_n = i_p = 0.$$

Circuit analysis can then be performed through all methods while observing these constraints. Op amps are not very useful by themselves; instead we use them in circuits designed to use these constraints. They are always designed to operate with “negative feedback” since they would be useless otherwise. In this class, we will always assume this to be the case. It is these summing-point constraints that given op amp circuits a wide diversity of important applications.

18.4 Applying Summing-Point Constraints

Example. Suppose ground is connected to a 5V source with voltage V_p leading to the non-inverting input. A 2V voltage source is connected from the output vertex, leading to a voltage V_n into the inverting input. The output splits from leading to the 2V source to the positive node. Determine V_0 .

We know that $V_p = 5V$. Secondly, we know that $V_p = V_n = 5V$ (virtual short circuit). Thus, since $V_n - V_0 = 2V$, we solve to find that $V_0 = 3V$. Finally, the 1000Ω resistor has no effect (ideally). The attached resistance does not affect voltage (since this is a property of an ideal voltage source). The voltage source inside the op-amp is essentially a dependent voltage source.

Example. Suppose we are given a ground connection leading to non-inverting input $+$. A $2mA$ current source is directed to the inverting input. Before it reaches the inverting input at $-$, it splits off into another branch that passes a 1000Ω resistor. This path joins with the path leading from the op-amp to the positive node of V_0 , where the negative node leads to ground. Determine V_0 .

We note that $V_p = 0$, and $i_p = 0$. V_n is also equal to 0, and since we assume infinite input resistance, $i_n = 0$. Thus, the current is diverted by the op-amp input since $i_n = 0$ in a virtual open circuit. Thus, across the 1000Ω resistance, we have a voltage of

$$V_1 = (2mA)(1000\Omega) = 2V.$$

Thus, since $V_1 = V_n - V_0$, we have $V_0 = 0V - 2V = -2V$.

19 March 1, 2017

19.1 Op-Amp Circuits Cont'd

There are many interesting and useful circuits that can be made. The simplest of them are some basic amplifier configurations:

The **inverting amplifier** is a ground connection to the non-inverting input. There is a voltage source V_{in} from $-$ to $+$ leading to a resistance of R_1 that leads to node a . This then leads to the inverting input, with another branch at a leading to resistance R_2 which connects with the path leading from the op-amp. This leads to the positive terminal of V_0 , where the negative terminal is connected to ground. Let us demonstrate an easy way to analyze op-amp problems:

1. All of the interesting properties occur at the op-amp input terminals. This is almost always our starting point.
2. Node equations at the input terminals tend to greatly simplify the job.

3. Seldom are we required to write node equations at the op-amp outputs. This is usually taken care of by the second point above.

With a node equation at a , the summing-point constraints tell us that $i_n = i_p = 0$, and $V_a = 0$ since V_d across the input terminals is 0. Thus,

$$\frac{V_a - V_{in}}{R_1} + \frac{V_a - V_0}{R_2} + i_n = 0,$$

where $i_n = 0$. Since we already know $V_a = 0$, we find that

$$V_0 = -\frac{R_2}{R_1}V_{in}.$$

This may also be written as

$$Av = \frac{V_0}{V_{in}} = -\frac{R_2}{R_1},$$

where A is the closed-loop gain. We did not write an equation at node V_0 . Doing so, we obtain

$$\frac{V_0 - V_a}{R_2} + i = 0$$

where i is the unknown current flowing to the output end of the op-amp. If we need i , then this is the equation that we use. However, since we were not asked to determine i , we can omit this step.

The **non-inverting input** is similar to the inverting amplifier, except that V_{in} is moved to the non-inverting (+) terminal. The summing-point constraint tells us that $i_n = i_p = 0$, and $V_a = V_{in}$ since both the + and - terminals are at V_{in} . Thus, at node a , we have

$$\frac{V_a}{R_1} + \frac{V_a - V_0}{R_2} + i_n = 0,$$

where $i_n = 0$. Since $V_a = V_{in}$, we can solve this to find that

$$V_0 = \left(1 + \frac{R_2}{R_1}\right)V_{in},$$

which may be written as

$$Av = \frac{V_0}{V_{in}} = 1 + \frac{R_2}{R_1},$$

where Av is the closed-loop gain. Note the result is positive (non-inverting amplifier).

20 March 3, 2017

20.1 Op-Amp circuits Cont'd

Another useful inverting amplifier circuit consists of two voltage sources connected from ground from $-$ to $+$, with one being V_a and the other V_b . These encounter resistances of R_A and R_B respectively before connecting at node a . This leads to the inverting input with a current i_n , and also leads to a path that connects to a resistance R_f before connecting with the output of the op-amp. This leads to the positive terminal of V_0 , with the negative terminal connected to ground. A ground connection with current i_p is connected to the non-inverting input. As before, we can write a node equation at node a ,

$$\frac{V_a - V_A}{R_A} + \frac{V_a - V_B}{R_B} + i_n + \frac{V_a - V_0}{R_f} = 0,$$

where $i_n = 0$. The op-amp imposes $V_a = 0$, $i_n = i_p = 0$, so we find V_0 where

$$\frac{0 - V_A}{R_A} + \frac{0 - V_B}{R_B} + \frac{0 - V_0}{R_f} = 0,$$

so

$$V_0 = -\frac{R_f}{R_A}V_A - \frac{R_f}{R_B}V_B.$$

If we choose $R_A = R_B = R$, then we obtain

$$V_0 = -\frac{R_f}{R}(V_A + V_B).$$

This is an inverting **summing amplifier** (for instance, this is a part of an audio mixing circuit).

The **differential amplifier** is another important and very important configuration. Ground is connected to both V_1 and V_2 . V_1 connects to a resistor R_3 before reaching node B . This is connected to ground through a resistor R_4 and to the non-inverting input with i_p . V_2 is connected to resistance R_1 before reaching node A . Node A leads to the inverting input with i_n , and also branches to R_2 which reconnects with the output of the op-amp. This leads to the positive terminal of V_0 , with the negative terminal leading to ground. This amplifier combines both an inverting and non-inverting amplifier. From the node equations at A and B ,

$$\frac{V_A - V_2}{R_1} + \frac{V_A - V_0}{R_2} + i_n = 0,$$

$$\frac{V_B}{R_4} + \frac{V_B - V_1}{R_3} + i_p = 0,$$

where $i_n = i_p = 0$. From the node equation at B , we find

$$V_B = V_1 \left(\frac{R_4}{R_3 + R_4} \right),$$

which is the voltage divider equation! Since there is no current that can flow because $i_p = 0$ (virtual open circuit), then R_3 and R_4 are in series. Now, we substitute $V_A = V_B$, where V_B is shown above, we find that

$$V_0 = \left(\frac{R_1 + R_2}{R_1} \right) \left(\frac{R_4}{R_4 + R_3} \right) V_1 - \left(\frac{R_2}{R_1} \right) V_2.$$

If we let $R_4 = R_2$ and $R_3 = R_1$, then we obtain

$$V_0 = \frac{R_2}{R_1} (V_1 - V_2),$$

which is the amplified voltage difference.

21 March 6, 2017

21.1 Examples of Other Op-Amp Circuits

Example. Consider the a ground connection to the non-inverting input with $i_p = 0$. A ground connection to V_{in} from $-$ to $+$ passes through R_1 before reaching node a . This leads to the inverting input with i_n and to node b through a resistance of R_2 . At node b , we have a connection to ground through a resistance of R_3 , and a connection to the output of the op-amp through resistance R_4 . This is connected to the positive terminal of V_0 with the negative terminal connected to ground. Determine V_0 .

At node a , we have

$$\frac{V_a - V_{in}}{R_1} + i_n + \frac{V_a - V_b}{R_2} = 0,$$

where $i_n = 0$. Since we have $V_a = 0$, we have

$$-\frac{V_{in}}{R_1} - \frac{V_b}{R_2} = 0.$$

For node b , we have

$$\frac{V_b - V_a}{R_2} + \frac{V_b}{R_3} + \frac{V_b - V_0}{R_4} = 0.$$

Thus,

$$V_0 = R_4 \left(\frac{1}{R_2} + \frac{1}{R_3} + \frac{1}{R_4} \right) V_b.$$

Substituting the node equations, we get

$$V_0 = -\frac{R_4 R_2}{R_1} \left(\frac{1}{R_2} + \frac{1}{R_3} + \frac{1}{R_4} \right) V_{in}.$$

Example. Consider the differential amplifier with V_1 connected to a 1000Ω resistor before reaching node b , which is connected to ground through a 1000Ω resistor and leading to the non-inverting input. V_2 reaches node a through a 2000Ω resistor leading to the inverting input. Node a also branches through a 2000Ω resistor to join with the output of the op-amp. Solve for V_0 .

First, we start with what we already know for certain. Since $i_p = 0$, the resistors are in series, so this is a voltage divider at the positive input terminal. Thus, $V_a = V_b = V_1/2$. At node a , we have

$$\frac{V_a - V_2}{2000} + \frac{V_a - V_0}{2000} + i_n = 0,$$

where $i_n = 0$. Thus, $V_0 = 2V_a - V_2$. Substituting the equations, we obtain $V_0 = V_1 - V_2$.

Example. Suppose we have a ground connection to V_{in} from $-$ to $+$ leading to a 4000Ω resistance to node a . Node a connects to the inverting input terminal and also through a 500Ω resistor to node b . Node b is connected from the output of the op-amp, and is also connected to node c through a 1000Ω resistor. Node c connects to the non-inverting input and to ground connected to V_{in} through a 4000Ω resistor. Find V_a .

As always, we first consider what happens at the input terminals. At node a ,

$$\frac{V_a - V_{in}}{4000} + i_n + \frac{V_a - V_b}{500} = 0,$$

where $i_n = 0$. From the summing point constraints, we know that $V_a = V_c$. Thus, we have two unknowns. At node c , we have

$$\frac{V_c - V_b}{1000} + i_p + \frac{V_c}{4000} = 0,$$

where $i_p = 0$. We solve these equations to get $V_a = -V_{in}$. Notice that we did not have to write an equation at node b . This op-amp circuit is emulating a negative resistance! Consider that $V_a = -V_{in}$, so

$$i = \frac{V_{in} - (-V_{in})}{4000} = \frac{2V_{in}}{4000}.$$

Thus, resistance is

$$R = \frac{V_a}{i} = -\frac{V_{in}}{i} = -\frac{V_{in}}{\frac{2V_{in}}{4000}} = -2000\Omega.$$

Example. Consider the following two op-amp problem. We are presented a $5V$ source from $-$ to $+$ from ground through a $20k\Omega$ resistor to node d . This is also

connected to a $10k\Omega$ resistance from ground, and also leads to the non-inverting input. The output leads to V_0 and also through a $100k\Omega$ resistor to node c . The inverting input connection is attached to node a . Node c is attached to the inverting input of the second op-amp, and is also connected to node b through a $40k\Omega$ resistor. Node b is also connected to the output of the second op-amp and leads to a through a $10k\Omega$ resistor. Node a is connected through a $20k\Omega$ resistor to a $1V$ source from $+$ to $-$ leading to ground. The non-inverting input of the second op-amp is connected to ground. Determine V_0 .

On the first op-amp, $i_{n1} = 0$, so we have a simple voltage divider. Thus

$$V_d = \frac{10}{10 + 20} * 5 = \frac{5}{3}V.$$

We note that this must also be the voltage at a . Additionally, we know that $V_c = 0V$ since it is the inverting terminal while the non-inverting terminal of the second op-amp has $0V$. We now solve a node equation at node a ,

$$\frac{V_a - 1}{20} + i_{n1} + \frac{V_a - V_b}{10} = 0,$$

where $i_{n1} = 0$. Since we know that $V_a = \frac{5}{3}V$, we determine that $V_b = 2V$. We consider the last point of interest of the input terminals at node c ,

$$\frac{V_c - V_0}{100} + \frac{V_c - V_b}{40} + i_{n2} = 0,$$

where $i_{n2} = 0$. With $V_b = 2V$ and $V_c = 0V$, we find that $V_0 = -5V$.

21.2 Input Resistance of Op-Amp Circuits

Consider an inverting amplifier. What want to determine the resistance seen by the source V_{in} . We recall that $V_a = 0$ due to the virtual short circuit. The current through the source and resistance is

$$i_1 = \frac{V_{in}}{R_1},$$

so the source sees a resistance of R_1 . Now, consider the non-inverting amplifier. We have $i_p = 0$, so

$$R_{in} = \frac{V_{in}}{i_p} = \infty.$$

But this is an open circuit!

21.3 Another Application of Op-Amps - Comparators

We have shown that for an inverting amplifier,

$$V_0 = -\frac{R_2}{R_1}V_{in}$$

with a closed loop gain of

$$Av = -\frac{R_2}{R_1}.$$

Resistor R_2 is a critical component here, since it provides the required negative feedback to allow this circuit to operate. Also recall that the op-amp itself requires an external power source to operate. Thus, with a ground connection to both $-15V$ and $15V$ from $-$ to $+$, both leading to opposite sides of an op-amp, these external power sources ensure that V_0 can be any voltage in the range of $-15V \leq V_0 \leq 15V$. Now, suppose that R_2 is not there. We are now running the op-amp open loop. Using the ideal model, this becomes a ground connection to V_i to a resistance leading to the negative terminal, and a ground connection leading to the positive terminal of V_d . This is expressed as a ground connection to a dependent voltage source from $-$ to $+$ of AV_d leading to the positive terminal, with a ground connection to the negative terminal of V_0 . Since we have $V_d = -V_{in}$, then

$$V_0 = -Av_{in},$$

where A is the open loop gain which is a large number (ideally infinite). We have two cases:

$$V_{in} < 0, V_0 = +[\text{huge number}],$$

$$V_{in} > 0, V_0 = -[\text{huge number}].$$

Here, the [huge number] is limited by the external power source, so

$$V_{in} < 0, V_0 = +15V,$$

$$V_{in} > 0, V_0 = -15V.$$

This behaviour makes the op-amp a very useful **voltage comparator**.

22 March 8, 2017

22.1 Inductors and Capacitors

So far, we have considered the basic circuit elements of sources and resistors. Inductors and capacitors are dependent on electromagnetic fields. **Capacitors** are from the separation of charge that produces an electric field. **Inductors** are from the

motion of charge that produces a magnetic field. Unlike resistors, these devices can store energy and return the stored energy (they are not producers of energy).

The **capacitor** is represented by the circuit symbol of two lines placed perpendicular to the circuit separating the conductor, optionally with a C . Like all circuit elements, the capacitor has its own important voltage-current relationship. For a current passing through the capacitor with a voltage from $+$ to $-$, we have

$$i = C \frac{dV}{dt},$$

where V is the voltage in Volts (V), i is the current in Amps (A), t is the time in seconds (s), and C is the capacitance in Farads (F).

The capacitor has constant voltage across the terminals that results in zero current flow. That is, the current looks like an open circuit. That is,

$$\frac{dV}{dt} = 0$$

when V is constant in the above expression. Furthermore, voltage cannot change instantaneously, since current would be infinite. We can manipulate the capacitance expression to find capacitor voltage in terms of current:

$$V(t) = \frac{1}{C} \int_{t_0}^t i(t) dt + V(t_0),$$

where it is usually assumed that $t_0 = 0$.

Example. *Plot voltage over time across a capacitor C through which a current i is flowing from $+$ to $-$ across the circuit element.*

The voltage over time $V(t)$ would increase from $V(0)$ as time increases, with the slope equal to $\frac{i}{C}$.

22.2 Power and Energy in the Capacitor

We again use the passive reference convention. When current is in the same direction as a voltage drop, $P = Vi$. This is then expressed as

$$P = V \left(C \frac{dV}{dt} \right),$$

or as

$$P = i \left(\frac{1}{C} \int_0^t i(t) dt + V(0) \right).$$

For energy, we recall that $P = \frac{dW}{dt}$, so solving with integration, we obtain

$$W = \frac{CV^2}{2},$$

where W is energy in Joules.

22.3 Capacitances in Series and Parallel

Applying KCL, we note that

$$\begin{aligned}
 i &= i_1 + i_2 + i_3 \\
 &= C_1 \frac{dV}{dt} + C_2 \frac{dV}{dt} + C_3 \frac{dV}{dt} \\
 &= (C_1 + C_2 + C_3) \frac{dV}{dt} \\
 &= C_{eq} \frac{dV}{dt}
 \end{aligned}$$

Thus, capacitances in parallel add since

$$C_{eq} = C_1 + C_2 + C_3.$$

23 March 10, 2017

23.1 Capacitances in Series and Parallel Cont'd

In series, we apply KVL to obtain

$$\begin{aligned}
 V &= V_1 + V_2 + V_3 \\
 &= \frac{1}{C_1} \int_0^t i(x) dx + \frac{1}{C_2} \int_0^t i(x) dx + \frac{1}{C_3} \int_0^t i(x) dx \\
 &= \frac{1}{C_{eq}} \int_0^t i(x) dx
 \end{aligned}$$

Thus, capacitances in parallel are like parallel resistors since

$$\frac{1}{C_{eq}} = \frac{1}{C_1} + \frac{1}{C_2} + \frac{1}{C_3}.$$

Example. Given a capacitor, find the voltage for a given current waveform. Assume that there is no initial charge on the capacitor. A current $i(t)$ passes a capacitor with voltage $V(t)$ from $+$ to $-$ where $C = 500\mu F = 5 * 10^{-4} F$. $i(t)$ is $20mA$ from 0 to 2 , is $0mA$ from 2 to 3 , and is $-20mA$ from 3 to 5 .

For a capacitor, we recall that

$$i(t) = C \frac{dV(t)}{dt},$$

and

$$V(t) = \frac{1}{C} \int_0^t i(x) dx + V(0),$$

where $V(0) = 0$. For the time $0 \leq t \leq 2$,

$$\begin{aligned} V(t) &= \frac{1}{5 * 10^{-4}} \int_0^t (20 * 10^{-3}) dx \\ &= \frac{20 * 10^{-3} x}{5 * 10^{-4}} \Big|_0^t \\ &= 40t \end{aligned}$$

For the time $2 < t \leq 3$,

$$\begin{aligned} V(t) &= \frac{1}{C} \int_2^t i(x) dx + V(2) \\ &= \frac{1}{C} \int_2^t 0 dx + V(2) \\ &= V(2) \\ &= 80V \end{aligned}$$

For the time $3 < t \leq 5$,

$$\begin{aligned} V(t) &= \frac{1}{C} \int_3^t i(x) dx + V(3) \\ &= -\frac{20 * 10^{-3} x}{5 * 10^{-4}} \Big|_3^t + 80 \\ &= -40(t - 3) + 80 \\ &= -40t + 200 \end{aligned}$$

We can then sketch this on a graph of $V(t)$ and t where the slope from 0 to 2 is 40, the slope from 2 to 3 is 0, and the slope from 3 to 5 is -40. At the last interval for $t > 5$, we have $i(t) = 0$, so $V(t) = V(5) = 0$. Now, with the waveform for current and voltage, we can determine the waveform for power by multiplying $P = Vi$. Thus, the capacitor begins absorbing energy before giving it back. We recall that energy

$$W = \frac{1}{2} CV^2.$$

23.2 Inductors

The circuit symbol for the inductor is curls within the conductor with the circuit parameter for inductance L underneath. The voltage-current relationship for an inductor L with a voltage V through which a current flows from + to - is given by

$$V = L \frac{di}{dt},$$

where V is the voltage in Volts (V), i is the current in Amperes (A), t is the time in seconds (s), and L is the inductance in Henrys (H).

For an inductor, a constant value of current causes zero voltage drop, so the inductor behaves like a short circuit. This is because

$$\frac{di}{dt} = 0$$

when i is constant. Furthermore, current cannot change instantaneously, since this would produce infinite voltage. We can manipulate the formula to find the current through the inductor in terms of voltage,

$$i(t) = \frac{1}{L} \int_{t_0}^t V(t) dt + i(t_0).$$

As before, we usually have $t_0 = 0$.

23.3 Power and Energy in the Inductor

When current is in the same direction as a voltage drop, $P = Vi$. This is then expressed as

$$P = Li \frac{di}{dt}.$$

For energy, we recall that $P = \frac{dW}{dt}$, so solving with integration, we obtain

$$W = \frac{Li^2}{2},$$

where W is energy in Joules.

23.4 Inductors in Series and Parallel

Applying KCL, we note that In series, we apply KVL to obtain

$$\begin{aligned} V &= V_1 + V_2 + V_3 \\ &= L_1 \frac{di}{dt} + L_2 \frac{di}{dt} + L_3 \frac{di}{dt} \\ &= (L_1 + L_2 + L_3) \frac{di}{dt} \end{aligned}$$

Thus, inductances in series add since

$$L_{eq} = L_1 + L_2 + L_3.$$

In parallel, we apply KCL to obtain

$$\begin{aligned} i &= i_1 + i_2 + i_3 \\ &= \frac{1}{L_1} \int_0^t V dx + \frac{1}{L_2} \int_0^t V dx + \frac{1}{L_3} \int_0^t V dx \\ &= \left(\frac{1}{L_1} + \frac{1}{L_2} + \frac{1}{L_3} \right) \int_0^t V dx \end{aligned}$$

Thus, inductances in parallel are like parallel resistances since

$$\frac{1}{L_{eq}} = \frac{1}{L_1} + \frac{1}{L_2} + \frac{1}{L_3}.$$

Example. Let an inductor of $2H$ have a voltage $V(t)$ with a current $i(t)$ flowing from $+$ to $-$. The waveform of $i(t)$ to t starts from 0 and reaches 3 at $0.1s$, then reaches 0 at 0.2 , -3 at 0.3 , and 0 at 0.4 . Given this inductor with $i(t)$, find $V(t)$, $P(t)$, and $W(t)$.

For the inductor, we recall that

$$V(t) = L \frac{di(t)}{dt}.$$

For the interval $0 < t \leq 0.1$,

$$\begin{aligned} V(t) &= L \frac{di(t)}{dt} \\ &= 2 \left(\frac{3}{0.1} \right) \\ &= 60V \end{aligned}$$

For the interval $0.1 < t \leq 0.3$,

$$\begin{aligned} V(t) &= L \frac{di(t)}{dt} \\ &= 2 \left(\frac{-6}{0.2} \right) \\ &= -60V \end{aligned}$$

For the interval $0.3 < t \leq 0.4$,

$$\begin{aligned} V(t) &= L \frac{di(t)}{dt} \\ &= 2 \left(\frac{3}{0.1} \right) \\ &= 60V \end{aligned}$$

We now sketch the $V(t)$ and $i(t)$ graph. $P = Vi$, so we multiply both graphs to obtain the $P(t)$ and t graph. For the $W(t)$ and t graph, we recall that $W = Li^2/2$. Plotting this using the graphs, we note that it is never negative, meaning that it is never producing energy.

23.5 Steady State Sinusoidal Analysis

So far, we have considered circuits in which sources are DC. We now investigate circuits where sources deliver sinusoidal (AC) currents and voltages. The methods of analysis are identical, but the arithmetic changes from real to complex.

23.6 Sinusoidal Currents and Voltages

Let $V(t) = V_m \cos(\omega t + \theta)$, where V_m is the peak value, ω is the angular frequency measured in radians/sec, and θ is the phase angle measured in radians. This can be plotted on a $V(t)$ and t graph, where V_m is the highest point, and θ is the angle which the maximum height at V_m is displaced on the x axis towards the right (the graph is shifted to the right by θ). The sinusoid is periodic with period T . We have one complete period when the angle increases by 2π . We note then that

$$\begin{aligned}\omega t|_{t=T} &= 2\pi \\ \omega T &= 2\pi \\ T &= \frac{2\pi}{\omega}\end{aligned}$$

Frequency is defined as the number of complete periods (cycles) per second, so

$$f = \frac{1}{T},$$

where f is the frequency in Hertz (Hz). We also have

$$\omega = \frac{2\pi}{T} \implies \omega = 2\pi f,$$

in radians/second. By convention, we use cosine and not sine. They are related by

$$\begin{aligned}\sin(\omega t) &= \cos(\omega t - \pi/2) \\ &= \cos(\omega t - 90^\circ)\end{aligned}$$

We say that $\sin(\omega t)$ has a phase angle of -90° .

23.7 Root-Mean-Square Values

We often express voltages and current in terms of their peak values V_m and i_m , but also in terms of their **root-mean-square (rms)** values. Consider power in a

resistor over one period of the waveforms. Instantaneous power is

$$P(t) = V(t)i(t) = \frac{V^2(t)}{R}.$$

The energy over one period is

$$E_T = \int_0^T P(t)dt.$$

An important measure is the **average power** over one period,

$$\begin{aligned} P_{avg} &= \frac{E_T}{T} \\ &= \frac{1}{T} \int_0^T P(t)dt \\ &= \frac{1}{T} \int_0^T \frac{V^2(t)}{R} dt \end{aligned}$$

This can be expressed as

$$P_{avg} = \frac{\left(\sqrt{\frac{1}{T} \int_0^T V^2(t)dt} \right)^2}{R} = \frac{V_{rms}^2}{R},$$

where the square root is the “root”, the V^2 is the square, and the means is the term under the radical. RMS values are sometimes called **effective values**. In the real world, AC voltages are specified in rms, not peak (for instance, household voltages are 120V). Power is also the **average power**, not instantaneous power (for instance, a 100W light bulb uses 100W of average power).

23.8 Relating to DC Circuits

$$V(t) = V_m \cos(\omega t + \theta),$$

where $\omega = 0$, $\theta = 0$, so

$$V(t) = V_m.$$

Also,

$$V(t) = V_{rms} = V_m,$$

$$i(t) = I_{rms} = I_m,$$

$$P(t) = P_{avg}.$$

For sinusoidal voltages and currents, peaks and rms values are not equal. It can be shown that the **Sinusoidal RMS Value** is

$$V_{rms} = \frac{V_m}{\sqrt{2}}.$$

The voltage in your home is $V_{rms} = 120V$. Since $\omega = 2\pi f = 2\pi * 60Hz$, this means that

$$\begin{aligned} V(t) &= 120\sqrt{2} \cos(\omega t + \theta) \\ &= 169.7 \cos(120\pi t + \theta) \end{aligned}$$

Example. Let $V(t) = 10 \sin(1000\pi t + 30^\circ)$. Express this as a cosine, give angular frequency, frequency in Hz, the rms voltage, and the average power in a 10Ω resistor.

We have

$$\begin{aligned} V(t) &= 10 \sin(1000\pi t + 30^\circ) \\ &= 10 \cos(1000\pi t + 30^\circ + 90^\circ) \end{aligned}$$

Thus, the angular frequency is $\omega = 1000\pi$ radians/second, and the frequency in Hertz is $f = \omega/(2\pi) = 500Hz$. The rms voltage is $V_{rms} = V_m/\sqrt{2} = 10V/\sqrt{2} = 7.071V$. The average power is therefore $P_{avg} = V_{rms}^2/R = 50V/10\Omega = 5W$. We can also sketch the instantaneous power where

$$P(t) = \frac{V^2(t)}{R} = \frac{100}{10} \cos^2(1000\pi t - 60^\circ).$$

We can now use the identity $\cos^2(x) = \frac{1}{2}(1 + \cos(2x))$ to rewrite this as

$$P(t) = 5 + 5 \cos(2000\pi t - 120^\circ).$$

This can then be sketched on a $P(t)$ and t graph where P_{avg} is shown by a translation of 5 units up, and 10 is the maximum height of the waveform.

24 March 15, 2017

24.1 Phasors

When dealing with sinusoidal voltages and currents, we need a convenient way to add them to satisfy KCL and KVL. Consider a circuit consisting of a voltage source $V(t)$ from $-$ to $+$ in the direction of current $i(t)$. It encounters circuit elements with voltages $V_1(t)$, $V_2(t)$, and $V_3(t)$ respectively from $+$ to $-$. Let $V_1(t) = 10 \cos(\omega t)$, $V_2(t) = 5 \cos(\omega t - 30^\circ)$, and $V_3(t) = 5 \cos(\omega t + 90^\circ)$. Find $V(t) = V_m \cos(\omega t + \theta)$. KVL must be satisfied by this circuit over all time, so

$$-V(t) + V_1(t) + V_2(t) + V_3(t) = 0$$

This means that

$$V(t) = 10 \cos(\omega t) + 5 \cos(\omega t - 30^\circ) + 5 \cos(\omega t + 90^\circ).$$

We need to manipulate this to the correct form. To accomplish this, we instead express voltages and currents in terms of **phasors**. Let $V_1(t) = V_1 \cos(\omega t + \theta_1)$. We note that ω is usually fixed in value throughout the circuit analysis problem. We have a pair of independent parameters describing the voltage, where V_1 is the magnitude (amplitude) and θ_1 is the phase angle. The basic idea behind phasors is that we represent this as a vector on a plane, then add the vector lengths. Consider the following example of phasor representation:

$V_a(t) = V_a \cos(\omega t + \theta_a)$ implies that \bar{V}_a has a magnitude of V_a and a phase angle of θ_a . $V_b(t) = V_b \sin(\omega t + \theta_b)$ is the same as $V_b \cos(\omega t + \theta_b - 90^\circ)$. This implies that \bar{V}_b has a magnitude of V_b and a phase angle of $\theta_b - 90^\circ$. Similarly, for a current of $i_c(t) = I_c \cos(\omega t + \theta_c)$, this implies that \bar{I}_c has a magnitude of I_c and a phase angle of θ_c .

24.2 Complex Numbers Review

Remark. The following section is just a basic review of complex numbers. This section is included for completion only.

To manipulate phasors, we need to make use of complex numbers. Complex numbers involve imaginary numbers. Since i represents currents, we shall denote the imaginary number as

$$j = \sqrt{-1}.$$

For a complex number, it is composed of a real part and an imaginary part. This can be plotted on a graph where the y axis represents the imaginary part and the x axis represents the real part of the number. Our number $x = 2 + j4$ for instance, would be represented as a point that is 2 units on the real axis and 4 units on the imaginary axis. x is therefore a point on the complex plane. The **complex conjugate** of a complex number consists of the real part of the original number along with the negative of the imaginary part. That is, the sign on the imaginary part is flipped. We can convert between the rectangular and polar forms of complex numbers. The rectangular form is what has been shown, while the polar form consists of a scalar M and an angle θ that the complex number forms on the complex plane.

To convert between polar in terms of M and θ , and rectangular in terms of $a + jb$, we note that,

$$\begin{aligned} M &= \sqrt{a^2 + b^2}, \\ \theta &= \tan^{-1} \left(\frac{b}{a} \right), \\ a &= M \cos(\theta), \\ b &= M \sin(\theta). \end{aligned}$$

Complex arithmetic must be done in rectangular form. Multiplication and division can be done in either form. In rectangular form, multiplication can be performed

as one would normally with variables. It is important to remember that $j^2 = -1$. In division, we multiply the numerator and denominator by the complex conjugate in order to simplify. In polar form, we multiply by multiplying the magnitudes $M = M_1 \cdot M_2$, and adding the angles $\theta = \theta_1 + \theta_2$. To divide, we divide the magnitudes $M = M_1/M_2$ and subtract the angles $\theta = \theta_1 - \theta_2$.

The key to phasors is the use of **Euler's identity**,

$$e^{j\theta} = \cos(\theta) + j \sin(\theta).$$

Multiplying both sides by M , we obtain

$$Me^{j\theta} = M \cos(\theta) + jM \sin(\theta),$$

where the left side is the complex exponential (another way of expressing M with θ), and the right side is the rectangular form.

Example. Determine the polar and rectangular forms of $x = 10e^{j30^\circ}$.

We note that the polar form is simply a magnitude of M with an angle of 30° . To compute the rectangular form, we use Euler's identity,

$$\begin{aligned} x &= 10e^{j30^\circ} \\ &= 10 \cos(30^\circ) + 10j \sin(30^\circ) \\ &= 8.66 + 5j \end{aligned}$$

24.3 KVL and KCL Using Phasors

Our reason for using phasors in this way is because it is simpler to use complex exponentials than trigonometric identities. The key is to express cosines as complex exponentials using Euler's identity,

$$\cos(x) = \operatorname{Re}(e^{jx}).$$

In our original KVL example, we had

$$V(t) = 10 \cos(\omega t) + 5 \cos(\omega t - 30^\circ) + 5 \cos(\omega t + 90^\circ).$$

Now, we can equivalently express this as

$$\begin{aligned}
 V(t) &= V_1(t) + V_2(t) + V_3(t) \\
 &= \operatorname{Re}(10e^{j\omega t}) + \operatorname{Re}(5e^{j(\omega t - 30^\circ)}) + \operatorname{Re}(5e^{j(\omega t + 90^\circ)}) \\
 &= \operatorname{Re}(10e^{j\omega t}) + \operatorname{Re}(5e^{j\omega t}e^{-j30^\circ}) + \operatorname{Re}(5e^{j\omega t}e^{j90^\circ}) \\
 &= \operatorname{Re}(10e^{j\omega t} + 5e^{j\omega t}e^{-j30^\circ} + 5e^{j\omega t}e^{j90^\circ}) \\
 &= \operatorname{Re}\left(\left(10 + 5e^{-j30^\circ} + 5e^{j90^\circ}\right)e^{j\omega t}\right) \\
 &= \operatorname{Re}\left(\left(10 + (4.33 - 2.5j) + 5j\right)e^{j\omega t}\right) \\
 &= \operatorname{Re}\left(\left(14.33 + 2.5j\right)e^{j\omega t}\right) \\
 &= \operatorname{Re}\left(\left(14.54e^{j9.90^\circ}\right)e^{j\omega t}\right) \\
 &= \operatorname{Re}\left(14.54e^{j(\omega t + 9.90^\circ)}\right) \\
 &= 14.54 \cos(\omega t + 9.90^\circ)
 \end{aligned}$$

We have added the three complex constant (phasors) inside the brackets separate from $e^{j\omega t}$ in Step 5 to determine the voltage.

24.4 Summary of Phasor Summation Method

First, we express the cosine functions as phasors. We then add the phasors. Afterwards, we convert the result back into a cosine function.

Example. Determine $V(t) = V_1(t) + V_2(t)$ and draw the phasor diagram when $V_1(t) = 5 \sin(\omega t + 45^\circ)$ and $V_2(t) = 10 \cos(\omega t + 90^\circ)$.

The “Time-domain” representation (cosine function) of $V_1(t) = 5 \cos(\omega t + 45^\circ - 90^\circ) = 5 \cos(\omega t - 45^\circ)$, while the time domain representation of $V_2(t)$ is as given. In phasor notation, this becomes $M_1 = 5$, $\theta_1 = -45^\circ$ and $M_2 = 10$, $\theta_2 = 90^\circ$ respectively. $\bar{V} = \bar{V}_1 + \bar{V}_2$. Thus,

$$\begin{aligned}
 \bar{V}_1 &= 5 \cos(-45^\circ) + j5 \sin(-45^\circ) \\
 &= 3.54 - 3.54j
 \end{aligned}$$

$$\begin{aligned}
 \bar{V}_2 &= 10 \cos(90^\circ) + j10 \sin(90^\circ) \\
 &= 0 + 10j
 \end{aligned}$$

Thus, converting back to polar, we have $\bar{V} = 3.54 + (10 - 3.54)j = 3.54 + 6.46j$, so we obtain

$$M = \sqrt{3.54^2 + 6.46^2} = 7.37.$$

$$\theta = \tan^{-1} \left(\frac{6.46}{3.54} \right) = 61.28^\circ.$$

Convertint back to the time-domain expression,

$$V(t) = 7.37 \cos(\omega t + 61.28^\circ).$$

A common misconception is that $\bar{V} = 3.54 + 6.46j$ means that the voltage is a complex number. However, the voltage is not a complex number. The voltage is a real-valued cosine with a magnitude of 7.37 and a plane angle of 61.28° .

25 March 20, 2017

25.1 Phase Relationship Between Sinusoids

Let $V_1(t) = 10 \cos(\omega t + 45^\circ)$, so \bar{V}_1 composed of M_1 and θ_1 is $M_1 = 10$ and $\theta_1 = 45^\circ$, and let $V_2(t) = 8 \cos(\omega t - 45^\circ)$, so \bar{V}_2 composed of M_2 and θ_2 is $M_2 = 8$ and $\theta_2 = -45^\circ$. We say that $V_1(t)$ is 90° higher in phase than $V_2(t)$, and therefore $V_1(t)$ leads $V_2(t)$ by 90° . Likewise, $V_2(t)$ lags $V_1(t)$ by 90° . On a graph, we see this since the top of $V_1(t)$ occurs before $V_2(t)$ by a 90° separation on the horizontal.

Example. Consider the phasor diagram where \bar{V}_1 has a magnitude of 7 at 135° , \bar{V}_2 has a magnitude of 10 at 30° , and \bar{V}_3 has a magnitude of 8 at 90° . Let $f = 100\text{Hz}$. Express each phasor voltage in the time domain as $V_m \cos(\omega t + \theta)$.

We note that we always represent angles as the angular distance from the positive real axis. Since $\omega = 2\pi f$, with f known, we solve to find that $\omega = 200\pi$ radians per second. Thus,

$$V_1(t) = 7 \cos(200\pi t + 135^\circ),$$

$$V_2(t) = 10 \cos(200\pi t + 30^\circ),$$

$$V_3(t) = 8 \cos(200\pi t + 90^\circ).$$

We note that $V_1(t)$ leads $V_2(t)$ by 105° and leads V_3 by 45° . $V_3(t)$ leads $V_2(t)$ by 60° and lags $V_1(t)$ by 45° .

25.2 Complex Impedances

Now, we need to revisit the voltage-current relationship for resistors, capacitors, and inductors when voltages and currents are sinusoidal.

For the inductor, we have

$$V_L(t) = L \frac{di_L(t)}{dt}.$$

Since we have $i_L(t) = I_m \cos(\omega t + \theta)$, the voltage will be

$$\begin{aligned} V_L(t) &= LI_m \frac{d}{dt} [\cos(\omega t + \theta)] \\ &= -\omega LI_m \sin(\omega t + \theta) \\ &= -\omega LI_m \cos(\omega t + \theta - 90^\circ) \\ &= \omega LI_m \cos(\omega t + \theta + 90^\circ) \end{aligned}$$

Expressed in phasor form, $i_L(t) = I_m \cos(\omega t + \theta)$ becomes \bar{I}_L with $M = I_m$ at θ degrees. On the other hand, $V_L(t) = \omega LI_m \cos(\omega t + \theta + 90^\circ)$ becomes \bar{V}_L with $M = \omega LI_m$ at $\theta + 90^\circ$ degrees. Through manipulation of the equation by factoring out \bar{I}_L and using Euler's identity, we find that

$$\bar{V}_L = j\omega L \cdot \bar{I}_L,$$

where the term $j\omega L$ is the **inductor impedance** Z_L . Thus, the impedance of an inductor in Ohms is also equal to $M = \omega L$ at 90° . We now have a phasor equivalent of Ohm's law, since

$$\bar{V}_L = Z_L \bar{I}_L.$$

Note that the voltage leads the current by 90° . That is, \bar{V}_L is 90° higher in phase than \bar{I}_L .

For the capacitor, we have

$$i_C(t) = C \frac{dV_C(t)}{dt}.$$

By the same analysis, we may write

$$\bar{V}_C = Z_C \bar{I}_C,$$

where Z_C is the **capacitor impedance** given by

$$Z_C = \frac{1}{j\omega C},$$

which can also be expressed as $M = \frac{1}{\omega C}$ at -90° . Note that voltage lags current by 90° . That is, \bar{V}_C is 90° less in phase than \bar{I}_C .

For the resistor, there is nothing new, since

$$\bar{V}_R = R \bar{I}_R,$$

where R is a real-valued constant (resistance). Furthermore, \bar{V}_R and \bar{I}_R are exactly in phase.

25.3 Summary of Impedances

1. **Inductor:**

$$\bar{V}_L = Z_L \bar{I}_L,$$

$$Z_L = j\omega L.$$

2. **Capacitor:**

$$\bar{V}_C = Z_C \bar{I}_C,$$

$$Z_C = \frac{1}{j\omega C}.$$

3. **Resistor:**

$$\bar{V}_R = R \bar{I}_R,$$

$$Z_R = R.$$

25.4 Circuit Analysis With Phasors and Complex Impedances

KVL and KCL must always be satisfied, whether it is AC or DC. Thus, the voltages around the loop must equal 0, and the currents entering a node must equal 0. For sinusoidal AC circuits, we express KVL and KCL in terms of phasors. That is,

$$\bar{V}_1 + \bar{V}_2 + \bar{V}_3 = 0,$$

$$\bar{I}_1 + \bar{I}_2 + \bar{I}_3 = 0.$$

We also use the phasor representation of voltage-current relationships, so

$$\bar{V} = Z \bar{I},$$

where Z is the complex impedance of an inductor, capacitor, or resistor. The analysis procedure that we follow is

- Use phasor for voltages and currents.
- Use complex impedance Z .
- Perform circuit analysis as usual.

Example (Complex Voltage Divider). *A circuit consists of a voltage source $V(t)$ from $-$ to $+$, which passes a 10Ω resistor before reaching a $500\mu F$ capacitor and a $100mH$ inductor. The voltage across the inductor from $+$ to $-$ is $V_L(t)$. $V(t) = 10 \cos(100t)$. Find $V_L(t)$.*

Like resistors in series, impedances add. Thus, we need to find the impedances of everything using $\omega = 100$ radians per second. The voltage $V(t)$ is \bar{V} with a magnitude of $M = 10$ at 0° . We already know the impedance of the resistor, since it is the same. So, we solve for the impedances of the capacitor and inductor, noting that $\frac{1}{j} = -j$,

$$Z_C = \frac{1}{j\omega C} = \frac{1}{j(100)(500 \cdot 10^{-6})} = -20j\Omega,$$

$$Z_L = j\omega L = j(100)(0.1) = 10j\Omega.$$

Thus, since $\bar{V} = 10$, we have

$$\begin{aligned} \bar{I} &= \frac{\bar{V}}{Z_R + Z_C + Z_L} \\ &= \frac{10}{10 - 20j + 10j} \\ &= \frac{10}{10 - 10j} \\ &= \frac{10}{10 - 10j} \cdot \frac{10 + 10j}{10 + 10j} \\ &= \frac{100 + 100j}{100 - 100j + 100j - 100j^2} \\ &= \frac{100 + 100j}{200} \\ &= 0.5 + 0.5j \end{aligned}$$

Now, we can determine \bar{V}_L , since $\bar{V}_L = Z_L \bar{I} = 10j \cdot (0.5 + 0.5j) = -5 + 5j$. Expressed in polar form, V_L has $M = 5\sqrt{2}$ and $\theta = 135^\circ$. Thus,

$$V_L(t) = 5\sqrt{2} \cos(100t + 135^\circ).$$

26 March 22, 2017

26.1 Complex Impedance Examples

Example. Suppose a $10 \cos(50t)$ voltage source connected to ground from $-$ to $+$ is connected to a $1000\mu F$ capacitor that reaches node a . This connects to ground through a 10Ω resistor, and to a $400mH$ inductor to node b . Node b is connected to ground through a $5 \sin(50t)$ current source in the reverse direction. Determine $V_a(t)$ and $V_b(t)$.

In terms of complex impedances, we have

$$Z_C = \frac{1}{j\omega c} = \frac{1}{j(50)(1000 * 10^{-6})} = -20j\Omega,$$

$$Z_L = j\omega L = j(50)(400 * 10^{-3}) = 20j\Omega,$$

$$Z_R = 10\Omega.$$

Now, the voltage expressed in phasor form has $M = 10$ and $\theta = 0$. the current source in phasor form has $M = 5$ and $\theta = -90$. It is thus given by $-5j$. Now, we write node-voltage equations,

$$\frac{\bar{V}_a - 10}{-20j} + \frac{\bar{V}_a}{10} + \frac{\bar{V}_a - \bar{V}_b}{20j} = 0,$$

$$\frac{\bar{V}_b - \bar{V}_a}{20j} - (-5j) = 0.$$

Solving these two equations, we find that $\bar{V}_a = \frac{90}{-1+2j} = -18-36j$ and $\bar{V}_b = 82-36j$. The phasor equation of \bar{V}_a therefore has $M = \sqrt{18^2 + 36^2} = 40.25$ with an angle of $\theta = -(180^\circ - \tan^{-1}(\frac{36}{18})) = -(180^\circ - 63.4^\circ) = -116.6^\circ$. Similarly, the phasor equation of \bar{V}_b has $M = \sqrt{82^2 + 36^2} = 89.55$ and $\theta = \tan^{-1}(-\frac{36}{82}) = -23.7^\circ$. Writing this in the time-domain representation, we obtain

$$V_a(t) = 40.25 \cos(50t - 116.6^\circ),$$

$$V_b(t) = 89.55 \cos(50t - 23.7^\circ).$$

Example. Consider the circuit composed of a $0.01 \cos(10^4 t)$ current source that connects to a $1K\Omega$ resistor, an inductor with a resistance of $200j\Omega$, and a capacitor with a resistance of $-200j\Omega$, all in parallel. Find the phasor voltage \bar{V} and all phasor currents.

First note that \bar{I} has $M = 0.01$ at an angle of 0° . We could start by finding the total impedance Z_{eq} , starting the the parallel branches of Z_C and Z_L . We find that

$$Z_{C,L} = \frac{(200j) \cdot (-200j)}{200j - 200j} = \infty.$$

the inductive impedance cancels the capacitive impedance! This is called **resonance**. Thus, we write node equations at a ,

$$-0.01 + \frac{\bar{V}}{1000} + \frac{\bar{V}}{200j} + \frac{\bar{V}}{-200j} = 0.$$

Solving this, we find that $\bar{V} = 10V$. Then,

$$\bar{I}_R = \frac{10}{1000} = 0.01A,$$

$$\bar{I}_L \frac{10}{200j} = -0.05jA,$$

$$\bar{I}_C = \frac{10}{-200j} = 0.05jA,$$

where $\bar{I}_R + \bar{I}_L + \bar{I}_C = 0.01$. Note that \bar{I}_C and \bar{I}_L cancel.

27 March 24, 2017

27.1 Thevenin Equivalent AC Circuits

As for DC circuits, we can reduce an AC circuit to a Thevenin or Norton equivalent. This is accomplished in the same way as for DC circuits, with the difference being that R_t is replaced with Z_t .

Example. Let $V(t) = 100 \cos(50t + 45^\circ)$ and $i(t) = 5 \cos(50t)$. The circuit is composed of a voltage source connected to ground at node b . From node a , we encounter a 10Ω resistor before reaching the voltage source from $+$ to $-$. Other parallel paths to node b from a include a path with a 10Ω resistor, a path with a current source $i(t)$ in the reverse direction, and a path with a $0.1H$ inductor and a $2000\mu F$ capacitor. Find the Thevenin equivalent.

The phasor representation of the current source is $\bar{I} = 5$ at an angle of 0 , while the phasor representation of the voltage source is $\bar{V} = 10$ at an angle of 45° . In terms of phasors and complex impedances,

$$Z_L = j\omega L = j(50)(0.1) = 5j\Omega,$$

$$Z_C = \frac{1}{j\omega C} = \frac{1}{j(50)(2000 * 10^{-6})} = -10j\Omega.$$

Applying a single node equation at node a , we have

$$\frac{\bar{V}_A - 100(\cos(45^\circ) + j\sin(45^\circ))}{10} + \frac{\bar{V}_A}{5j - 10j} - 5 + \frac{\bar{V}_A}{10} = 0.$$

Solving this, we obtain

$$\bar{V}_A = \frac{120.7 + 70.1j}{2 + 2j}.$$

We solve this by converting to polar form to make the division easier. We find that $\bar{V}_t = \bar{V}_A = 49.4$ at -14.4° . Thus, $V(t) = 49.4 \cos(50t - 14.4^\circ)$. Now, we find Z_t . For this example, we have no dependent sources! Thus, we zero the independent sources. This leaves us with two 10Ω resistor and a $-5j\Omega$ resistor, all in parallel. Thus,

$$\begin{aligned} Z_t &= \left(\frac{1}{10} - \frac{1}{5j} + \frac{1}{10} \right)^{-1} \\ &= \left(\frac{1}{5} - \frac{1}{5j} \right)^{-1} \\ &= \left(\frac{-5 + 5j}{25j} \right)^{-1} \\ &= \frac{125 - 125j}{50} \\ &= \frac{5}{2} - \frac{5}{2}j\Omega \end{aligned}$$

Not that this is a resistor and a capacitor in series, since Z_t can be expressed as a 2.5Ω resistor and a capacitor with a reactance of $-2.5j\Omega$ (the negative indicates capacitance).

28 March 17, 2017

28.1 Frequency Dependent Circuits

The frequency-dependent nature of inductors and capacitors paves way for a wide number of applications of AC circuits. Consider the circuit consisting of a voltage source $V(t)$ from $-$ to $+$ that leads to an inductor L , and a resistor with voltage $V_R(t)$ from $+$ to $-$. Let $V(t)$ be as follows:

$$V(t) = V_m \cos(\omega t + 0^\circ),$$

where V_m is the constant value and ω is the angular frequency (kept as a variable). In terms of phasors and complex impedances, we replace $V(t)$ with V_m at an angle of 0 , L with the complex impedance $Z_L = j\omega L$, and R with \bar{V}_R . This is a simple voltage divider, since

$$\bar{V}_R = \frac{R}{R + j\omega L} \cdot V_m,$$

where Z_L is **frequency-dependent**. We may write this as

$$\bar{V}_R = \frac{1}{1 + j\omega \left(\frac{L}{R}\right)} \cdot V_m.$$

Here, the magnitude and phase of \bar{V}_R are frequency-dependent. This dependence is called **frequency response**.

In the time domain,

$$V_R(t) = V_R \cos(\omega t + \theta_R),$$

where the amplitude $\|V_R\|$ is at frequency ω and θ_R is the phase of $V_R(t)$ at ω . The amplitude is

$$\begin{aligned} \|\bar{V}_R\| &= V_R \\ &= \left| \frac{1}{1 + j\left(\frac{\omega L}{R}\right)} \cdot V_m \right| \\ &= \frac{V_m}{\left| 1 + j\omega \left(\frac{L}{R}\right) \right|} \\ &= \frac{V_m}{\sqrt{\left(\frac{\omega L}{R}\right)^2 + 1^2}} \end{aligned}$$

We can graph V_R against ω . Doing so, we find that when ω is 0, we have V_m . As ω increases however, V_R drops.

This is called a **lowpass filter**. Sinusoids with low frequencies come through strongly, while higher frequencies come through at a reduced amplitude. An ideal lowpass filter has the frequency response that resembles a square in that V_R is V_m until a certain cutoff frequency ω_c at which V_R drops to 0. This can be compared to our graph, where the L-R circuit causes a less drastic descent to 0 at increased ω . For the ideal filter,

$$V_R = \begin{cases} V_m, & 0 \leq \omega \leq \omega_c \\ 0, & \omega > \omega_c \end{cases}$$

We usually define the cutoff frequency as the frequency at which $V_R = \left(\frac{1}{\sqrt{2}}\right) \cdot V_m$. By equating this with the expression for our L-R circuit, we find that

$$\frac{\omega_c L}{R} = 1.$$

Thus, $\omega_c = R/L$ radians per second and $f_c = \omega_c/2\pi$.

There are many other extremely useful filter circuits. One such **bandstop** filter circuit that exploits the series cancellation of impedances is composed of \bar{V}_{in} from $-$ to $+$ connected to a resistance R . This leads to a node which is connected to a capacitor with $Z_C = 1/j\omega C = -j/\omega C$, and an inductor $Z_L = j\omega L$, where the voltage across both the capacitor and inductor from $+$ to $-$ is \bar{V}_{out} . This leads back to the voltage source. The nodes where \bar{V}_{out} ends and begins lead to two other nodes. Thus, we find that

$$\bar{V}_{out} = \frac{Z_L + Z_C}{Z_L + Z_C + Z_R} \cdot \bar{V}_{in},$$

where the numerator will cancel at the frequency where $Z_L = -Z_C$, so $\bar{V}_{out} = 0$. Plotting \bar{V}_{out} against ω , we note that at $\omega = 0$, we are at V_m . The graph then drops to 0 when $\omega L = 1/\omega C$ (resonance) before rising again with increasing ω . This is called a **notch filter** and has many important applications. For instance, the “hum” in an audio system is due to a $60Hz$ power source.

29 March 29, 2017

29.1 Superposition in AC Circuits

As with other methods of AC circuit analysis, the procedure is identical to that of DC circuits. However, we use complex algebra! This method is the only way to analyze circuits with sources of different frequencies.

29.2 Power in AC Circuits

Consider an arbitrary complex impedance consisting of a voltage source from $-$ to $+$ of V_m at an angle of 0 , through which a current of \bar{I} flows to a resistor R and a complex impedance of jX before returning to the negative terminal of the voltage source. The load resistance is

$$Z = R + jX,$$

so that

$$\begin{aligned}\|Z\| &= \sqrt{R^2 + X^2}, \\ \theta &= \tan^{-1} \left(\frac{X}{R} \right).\end{aligned}$$

In the above relations, R is the resistive part, while X is the reactive part. We now have current phasor \bar{I} , which is

$$\bar{I} = \frac{V_m \angle 0^\circ}{\|Z\| \angle \theta} = \frac{V_m}{\|Z\|} \angle -\theta.$$

Thus, we let I_m be the magnitude, so

$$\bar{I} = I_m \angle -\theta.$$

We will now investigate four cases:

1. Resistor
2. Inductor
3. Capacitor
4. General Load

In a **Purely Resistive Load**, $X = 0$. We have

$$V(t) = V_m \cos(\omega t),$$

$$i(t) = I_m \cos(\omega t).$$

Power is $V(t)i(t)$, so $P(t) = V_m I_m \cos^2(\omega t)$. Using the identity $\cos^2(x) = \frac{1}{2}(1 + \cos(2x))$, we obtain

$$P(t) = \frac{1}{2} V_m I_m (1 + \cos(2\omega t)).$$

We note that if we plot $P(t)$ against t , we will obtain a sinusoidal wave with the maximum at $V_m I_m$, the minimum at 0 , and $P_{avg} = \frac{1}{2} V_m I_m$. The graph is always positive, as the resistor only absorbs power.

In a **Purely Inductive Load**, $R = 0$ and $X > 0$. For the inductor, $Z = j\omega L = \omega L \angle 90^\circ$, so $\theta = 90^\circ$. Thus,

$$V(t) = V_m \cos(\omega t),$$

$$i(t) = I_m \cos(\omega t - 90^\circ) = I_m \sin(\omega t).$$

Power is $V_m I_m \cos(\omega t) \sin(\omega t)$, but we use the identity $\cos(x) \sin(x) = \frac{1}{2} \sin(2x)$ to obtain

$$P(t) = \frac{V_m I_m}{2} \sin(2\omega t).$$

Plotting $P(t)$ against t , we find that the maximum occurs at $\frac{V_m I_m}{2}$, and the minimum occurs at $-\frac{V_m I_m}{2}$. The positive region indicates that energy is being absorbed, while the negative region indicates that energy is being given back. This is called **reactive power** - the average is zero.

In a **Purely Capacitive Load**, $R = 0$ and $X < 0$. We have $Z = \frac{1}{j\omega C} = \frac{1}{\omega C} \angle -90^\circ$, so $\theta = -90^\circ$. Thus,

$$V(t) = V_m \cos(\omega t),$$

$$i(t) = I_m \cos(\omega t + 90^\circ) = -I_m \sin(\omega t).$$

Power is therefore

$$P(t) = -\frac{V_m I_m}{2} \sin(2\omega t).$$

This is also reactive power. For the capacitor and inductor, no average power is consumed or generated.

For a **General Load** where $R \neq 0$ and $X \neq 0$, we allow for both resistance and capacitance or inductance. Thus, we allow θ in the range

$$-90^\circ \leq \theta \leq 90^\circ,$$

where the lower limit of -90° is purely capacitive and the upper limit of 90° is purely inductive. We have

$$V(t) = V_m \cos(\omega t),$$

$$i(t) = I_m \cos(\omega t - \theta),$$

$$P(t) = V_m \cos(\omega t) I_m \cos(\omega t - \theta),$$

which we can manipulate to obtain

$$P(t) = \frac{V_m I_m}{2} \cos(\theta) (1 + \cos(2\omega t)) + \frac{V_m I_m}{2} \sin(\theta) \sin(2\omega t).$$

Since the average of $\cos(2\omega t)$ and $\frac{V_m I_m}{2} \sin(\theta) \sin(2\omega t)$ are both 0, this means that the average power is

$$P = P_{avg} = \frac{V_m I_m}{2} \cos(\theta).$$

This is the power in Watts, absorbed by the resistive component of the total impedance. We recall that $V_{rms} = \frac{V}{\sqrt{2}}$ and $I_{rms} = \frac{I}{\sqrt{2}}$ to obtain

$$P = V_{rms} I_{rms} \cos(\theta),$$

where P is the average, or real, power. Note that for a resistor when $\theta = 0^\circ$, the \cos term cancels out. This term is very important, and is referred to as the **power factor**, where

$$PF = \cos(\theta).$$

In the general case, the **power angle** is

$$\theta = \theta_V - \theta_I.$$

We often state the PF and specify whether the current leads or lags voltage.

30 March 31, 2017

30.1 Power in AC Circuits Example

Example. *A load has a leading power factor of 0.707. Determine whether this is capacitive or inductive, and the power angle.*

A leading PF means that the current is leading (has a higher phase than) the voltage. Reviewing what we know, we have $Z = R + Xj = \|Z\|/\underline{\theta}$. Let

$$\bar{V} = V_m/\underline{\theta}_V,$$

$$\bar{I} = I_m/\underline{\theta}_I,$$

where we are given $\theta_I > \theta_V$ since the current leads the voltage. We know that $\bar{I} = \bar{V}/Z$, so $Z = \bar{V}/\bar{I}$. Thus,

$$\begin{aligned} Z &= \frac{V_m/\underline{\theta}_V}{I_m/\underline{\theta}_I} \\ &= \|Z\|/\underline{\theta}_V - \theta_I \end{aligned}$$

We note that the power angle $\theta_V - \theta_I$ is negative because $\theta_I > \theta_V$. What we now know is that since the power factor $PF = \cos(\theta) = 0.707$, and $\theta < 0$, this means that the power angle $\theta = -45^\circ$. This suggests that $Z = R + Xj$, where $X < 0$. The load therefore has a capacitance of $Z_C = -j/\omega C$.

Average **reactive power** is always zero. However, its instantaneous value is sinusoidal with peak value Q , where it is given by

$$Q = V_{rms}I_{rms} \sin(\theta).$$

This is flowing back and forth between the inductors/capacitors and the source. This might be a problem in large scale systems. Power companies may penalize you for reactive power. The units for reactive power are **Volt-Amperes-Reactive, VARs**.

Apparent power is a measure of the total power (average and reactive), and is given by

$$P_{app} = V_{rms}I_{rms}.$$

The units of apparent power are **Volt-Amps, VAs**. For instance, a $5kW$ load, a $10kVA$ load and a $15kVAR$ load means that we have $P = 5000W$, $V_{rms}I_{rms} = 10000VA$, and $Q = 15000VAR$ respectively.

30.2 The Power Triangle and Other Power Relationships

Each of P , Q , and apparent power can be represented in a triangle. In a right triangle, P is the adjacent side, Q is the opposite side, and $V_{rms}I_{rms}$ is the hypotenuse. The angle θ between the adjacent and the hypotenuse is the power angle, where θ is positive if inductive, and negative if capacitive.

It is easy to calculate P , Q , and apparent power directly from impedance. We have $Z = \|Z\|\underline{\theta} = R + Xj$. Additionally, we recall that $\cos(\theta) = R/\|Z\|$ and $\sin(\theta) = X/\|Z\|$. We also have

$$P = \frac{V_m I_m}{2} \cos(\theta) = \frac{V_m I_m}{2} \cdot \frac{R}{\|Z\|},$$

$$I_m = \frac{V_m}{\|Z\|}.$$

Thus, substituting V_m , we find that $P = I_m^2 R/2$. By applying other known expressions for the variables, we obtain

$$P = I_{rms}^2 R,$$

$$Q = I_{rms}^2 X,$$

$$P_{app} = \sqrt{P^2 + Q^2},$$

where these expressions represent the average power in Z , the reactive power in Z , and the apparent power in Z respectively.

Finally, **complex power** is defined as

$$\bar{S} = \frac{1}{3} \bar{V} \bar{I}.$$

Thus, we can expand \bar{V} and \bar{I} to obtain

$$\begin{aligned} \bar{S} &= \frac{1}{2} (V_m \underline{\theta}_V) \cdot (I_m \underline{-\theta}_I) \\ &= \frac{1}{2} V_m I_m \underline{\theta}_V - \theta_I \end{aligned}$$

where $\theta = \theta_V - \theta_I$ is the power angle. Expanding \bar{S} into rectangular form, we obtain

$$\begin{aligned}\bar{S} &= \frac{V_m I_m}{2} \cos(\theta) + j \frac{V_m I_m}{2} \sin(\theta) \\ &= P + jQ\end{aligned}$$

The apparent power can then be seen to be the magnitude of the complex power, as

$$P_{app} = \|\bar{S}\| = \sqrt{P^2 + Q^2}.$$

Example. We have a circuit with a voltage source of 1414 at an angle of 30° from $-$ to $+$ through which current \bar{I} passes. This splits off into two parallel paths that reconnect with the $-$ end of the voltage source. The first path has current \bar{I}_A with $10kVA$, $PF = 0.5$, and is leading. The second path has a current \bar{I}_B with $5kW$, $PF = 0.7$, and is lagging. Find \bar{I} in the example below.

The first load is specified in terms of applied power P_{app} in kVA , while the second load is specified in terms of average power P in Watts. Applying the power triangle to branch A , we have the hypotenuse as $10kVA$, with a power factor $\cos(\theta_A) = 0.5$. We also know that it is leading. Recall that a leading power factor means that the current leads the voltage. Thus, $\theta_I > \theta_V$. Thus, $\theta_A = \theta_V - \theta_I$ has a negative angle. The power angle is therefore given by

$$\theta_A = -(\cos^{-1}(0.5)) = -60^\circ.$$

We can now calculate P_A and Q_A for branch A ,

$$\begin{aligned}P_A &= V_{rms} I_{rms} \cos(\theta_A) \\ &= 10000 \cdot 0.5 \\ &= 5000W\end{aligned}$$

$$\begin{aligned}Q_A &= V_{rms} I_{rms} \sin(\theta_A) \\ &= -10000 \cdot 0.866 \\ &= -8.66kVAR\end{aligned}$$

Analogously, we find that $\theta_B = \cos^{-1}(0.7) = 45.57^\circ$. Hence, with our knowledge of P_B , we can find Q_B . We note that $\tan(\theta_B) = Q_B/P_B$, so

$$\begin{aligned}Q_B &= P_B \tan(\theta_B) \\ &= 5000 \tan(45.57^\circ) \\ &= 5.101kVAR\end{aligned}$$

The total power in both loads is therefore

$$P = P_A + P_B = 5kW + 5kW = 10kW,$$

$$Q = Q_A + Q_B = -8.660kVAR + 5.101kVAR = -3.559kVAR,$$

$$\bar{S} = P + jQ = (10000 - j3559)VA.$$

In polar form, this is

$$\bar{S} = 10610\angle-19.59^\circ,$$

where the negative angle indicates that current leads voltage. We know the total complex power and voltage, so by applying $\bar{S} = \frac{1}{2}\bar{V}\bar{I}$, we find that

$$\begin{aligned}\bar{I} &= \frac{2\bar{S}}{\bar{V}} \\ &= \frac{2 \cdot 10610\angle-19.59^\circ}{1414\angle30^\circ} \\ &= 15.0\angle-49.59^\circ\end{aligned}$$

Remark. When the power factor is leading, $\theta_I > \theta_V$, so the power angle is negative. When the power factor is lagging, $\theta_I < \theta_V$, so the power angle is positive.

31 April 3, 2017

31.1 DC Motors

We will now study electric motors and generators. Motors convert electrical energy to mechanical energy, while generators do the reverse. This is accomplished through electromechanical conversion.

Motors (and generators) are constructed with two major components, the **stator** (stationary part) and the **rotor** (rotating part). The rotor is connected to a shaft that connects to a mechanical load. Depending on the machine type, the rotor and stator contain conductors wired in coils called **windings**. This produces interacting magnetic fields, thereby producing physical torque. We note that **torque** is the twisting force that tends to cause rotation. The stator produces a magnetic field. This is often produced by the stator's field windings, or a permanent magnet. Motors can be found in many places (it constitutes 2/3 of power consumed in North America):

- Fans and Ventilation.
- Vacuum Cleaners.
- Rock Crushers.
- Trains.
- Disk Drives and Robotic Systems.

31.2 Operating Characteristics of Motors

Efficiency is a very important motor parameter. We can consider a voltage source with both ends connected to a motor, providing electrical input power P_{in} . The motor rotates the shaft which produces mechanical output power P_{out} . Efficiency is defined as

$$\eta = \frac{P_{out}}{P_{in}} \cdot 100\%.$$

For a DC machine, $P_{in} = Vi$ in Watts. Mechanical power output is given by

$$P_{out} = T_{out}\omega_m,$$

where T_{out} is the output torque in $N \cdot m$, ω_m is the angular shaft speed in $rads/sec$, and P_{out} is in Watts. The angular shaft speed can be expressed as

$$\omega_m = n_m \cdot \frac{2\pi}{60},$$

where n_m is the shaft speed in revolutions per minute. Note that $1HP = 746W$.

31.3 Speed Regulation

Depending on the motor type, speed may decrease with load. Speed regulation SR is defined as

$$SR = \frac{n_{no-load} - n_{full-load}}{n_{full-load}} \cdot 100\%,$$

where a smaller value is preferred. Values greater than 100% are possible.

Example. Given a DC motor with a 50HP rating, we find from measurements at the motor that $V = 220V$, $n_{no-load} = 1200rpm$, and $n_{full-load} = 1150rpm$. Under a full (rated) load, the power loss is equal to 3350W. At full load, find the efficiency, speed regulation, and input current.

To find efficiency, the motor is delivery 50HP of power, so $P_{out} = 50 \cdot 746 = 37300W$. The total power delivered plus the total lost is $P_{total} = 37300 + 3350 = 40650W$. Thus, this is the total input power. Efficiency is therefore

$$\eta = \frac{P_{out}}{P_{in}} = \frac{37300}{40650} \cdot 100\% = 91.76\%.$$

The input current and speed regulation can also be found,

$$i = \frac{P_{in}}{V} = \frac{40650}{220} = 184.77A,$$

$$SR = \frac{n_{no-load} - n_{full-load}}{n_{full-load}} \cdot 100\% = \frac{1200 - 1150}{1150} \cdot 100\% = 4.35\%.$$

32 April 5, 2017

32.1 Electrical Circuit of DC Motors

DC motors can be modeled with two simple circuits. The **field** consists of a current I_F flowing across a field resistance R_F and field windings (inductor) L_F , with a voltage across the entire system being V_F from + to -. The **armature** consists of a current I_A flowing across an armature resistance R_A and the shaft with speed ω_m and torque T_m . The voltage across the shaft is E_A from + to -, while the voltage across the entire system including the resistance is V_T from + to -. For a rotating DC machine, we have ω_m as the rotational speed in radians per second, and T_m is the torque in Newton meters.

Since we are operating in DC, the field current reduces to simply an expression over the resistance (L_F acts as a short circuit, so it is not considered). The induced armature voltage is given by

$$E_A = K\phi\omega_m,$$

where ϕ is the magnetic flux, and K is a machine constant. The total developed mechanical torque is

$$T_{dev} = T_m = K\phi I_A.$$

The total developed mechanical power is

$$P_{dev} = T_{dev}\omega_m.$$

Together, these three equations are the key to analyzing DC motor and generator circuits.

$$E_A = K\phi\omega_m,$$

$$T_{dev} = K\phi I_A,$$

$$P_{dev} = T_{dev}\omega_m.$$

We normally consider K and ϕ together, where $K\phi$ is the machine constant.

32.2 Magnetization Curve

The magnetization curve plots E_A against I_F . It consists of a linear region where the two are linearly dependent, until it reaches a particular shaft speed n_m . We then reach magnetic core saturation, where the slope levels out to 0. This is a typical magnetization curve for a given speed. A point on this curve gives us $K\phi$. From this, we can calculate the other values. Note that we may not always obtain the same curve, but $K\phi$ can almost always be calculated from the information given.

Example. *We have a DC motor that obeys the above curve. Additionally, we have $n_m = 1500\text{rpm}$, $P_{dev} = 10\text{HP}$, $I_F = 3\text{A}$, $R_A = 0.3\Omega$, and $R_F = 50\Omega$. Determine the developed torque, the armature I_A , the applied voltage V_T , and the efficiency.*

We can immediately determine the following,

$$\begin{aligned}\omega_m m &= n_m \cdot \frac{2\pi}{60} = 157.1 \text{ rads/s}, \\ P_{dev} &= 10 \text{HP} \cdot 746 = 7460 \text{W}, \\ T_{dev} &= \frac{P_{dev}}{\omega_m} = 47.49 \text{Nm}.\end{aligned}$$

We recall that the armature circuit consists of V_T attached to $R_A = 0.3\Omega$ with a current I_A . This then reaches the motor with an induced armature voltage E_A , where $P_{dev} = E_A I_A = 7460 \text{W}$. The motor is where the electrical world meets the mechanical world. Making use of the equation above, we need E_A to find I_A . From the given magnetization curve we have

$$\begin{aligned}E_A &= 200 \text{V}, \\ I_F &= 3 \text{A}, \\ n_m &= 1200 \text{rpm}.\end{aligned}$$

making use of the machine equation for E_A , we can rearrange to obtain

$$K\phi = \frac{E_A}{\omega_m} = \frac{200 \text{V}}{1200 \cdot \frac{2\pi}{60}} = 1.59.$$

Our motor is run at $n_m = 1500 \text{rpm}$, so

$$E_A = K\phi\omega_m = 1.59 \cdot 1500 \cdot \frac{2\pi}{60} = 250 \text{V}.$$

We can now use the above equation to find

$$I_A = \frac{P_{dev}}{E_A} = \frac{7460 \text{W}}{250 \text{V}} = 29.84 \text{A}.$$

Alternatively, we note that we could have used the machine equation for T_{dev} to find I_A , since

$$I_A = \frac{T_{dev}}{K\phi} = \frac{47.49}{1.59} = 29.84 \text{A}.$$

The applied voltage V_T can then be found by applying KVL, since

$$-V_T + I_A R_A + E_A = 0.$$

Solving this gives $V_T = (29.84 \text{A})(0.3\Omega) + 250 \text{V} = 258.95 \text{V}$. We recall that the total developed power in the armature was $P_{dev} = 10 \text{HP} = 7460 \text{W}$. The total input power is given by the power supplied by V_T and the field losses. Thus,

$$P_{in} = V_T I_A + I_F^2 R_F = (258.95)(29.84) + (3)^2(50) = 8177.1 \text{W}.$$

We can now calculate efficiency,

$$\eta = \frac{P_{dev}}{P_{in}} \cdot 100\% = \frac{7460}{8177.1} \cdot 100\% = 91.2\%.$$

32.3 Power and Torque: Developed vs. Output

At the mechanical output of the motor, we have electrical power from $P + I_A E_A$, and the developed mechanical power $P_{dev} = P = I_A E_A$ and $T_{dev} = \frac{P_{dev}}{\omega_m} = K\phi I_A$. Developed power and torque do not take into account rotational losses such as friction (bearings) and windage (wind resistance). In a practical motor, we have rotational losses P_{rot} and T_{rot} . Thus,

$$P_{out} = P_{dev} - P_{rot},$$

$$T_{out} = T_{dev} - T_{rot}.$$

If there are no rotational losses, then $P_{out} = P_{dev}$ and $T_{out} = T_{dev}$.

33 April 7, 2017

33.1 Shunt-Connected DC Machines

Suppose we are given a circuit with voltage source V_T from $-$ to $+$ with current I_L flowing to node a . Here, the path splits off into a path with a current of I_F through resistors R_{adj} and R_F and inductor L_F . The other path has a current of I_A through a resistor R_A , which then leads to a motor with E_A from $+$ to $-$, ω_m and T_{dev} . These two paths join together and meet at the negative terminal of the voltage source.

In the above machine configuration, the field and armature circuits are connected in parallel. The variable resistor R_{adj} is denoted with a line through a normal resistor symbol, and is available to adjust the torque-speed characteristic. The total input power is

$$P_{in} = V_T I_L,$$

where I_L is the total **line current**. Some of this creates the field. Power that is absorbed by the field is dissipated as heat,

$$P_F = I_F^2 (R_{adj} + R_F) = \frac{V_T^2}{R_{adj} + R_F}.$$

The armature resistance similarly dissipates power as heat,

$$P_A = I_A^2 R_A = \frac{(V_T - E_A)^2}{R_A}.$$

The remaining power is developed power P_{dev} , where

$$P_{dev} + E_A I_A,$$

$$T_{dev} = \frac{P_{dev}}{\omega_m} = \frac{E_A I_A}{\omega_m}.$$

Example. Consider a shunt-connected DC machine with $V_T = 200V$, $R_F = 10\Omega$, $R_{adj} = 20\Omega$, and $R_A = 0.065\Omega$. This machine has rotational losses (friction) represented by constant torque $T_{rot} = 12Nm$ (rotational power loss is proportional to speed, $P_{rot} = T_{rot}\omega_m$). From power tests on this machine when $I_F = 10A$ and $n_m = 1200rpm$, $E_A = 300V$. Additionally, the total required torque by the mechanical load is $T_{out} = 200Nm$. Find the motor speed and efficiency.

In the field, since we are dealing with DC, we have

$$I_F = \frac{300}{20 + 10} = 10A.$$

From the information given at $n_m = 1200rpm$, we know $I_F = 10A$ and $E_A = 300V$. We use the basic machine equations. This gives

$$K\phi = \frac{E_A}{\omega_m} = \frac{300}{1200 \cdot \frac{2\pi}{60}} = 2.387.$$

The total torque required by the load is $T_{out} = 200Nm$. Adding the rotational losses, we find

$$T_{dev} = T_{out} + T_{rot} = 200 + 12 = 212Nm.$$

Our strategy now is to find I_A , E_A , and then the speed. Making use of the machine equation,

$$I_A = \frac{T_{dev}}{K\phi} = \frac{212}{2.387} = 88.8.$$

From KVL, we have

$$E_A = V_T - I_A R_A = 300 - (88.8)(0.065) = 294.2V.$$

Therefore,

$$\omega_m = \frac{E_A}{K\phi} = \frac{294.2}{2.387} = 123.6rad/s,$$

$$n_m = \omega_m \cdot \frac{60}{2\pi} = 1177rpm.$$

We can then find efficiency after finding the input and output power,

$$P_{out} = T_{out}\omega_m = (200)(123.6) = 24652W,$$

$$P_{in} = V_T I_L = 300(I_F + I_A) = 300(10 + 88.8) = 29640W,$$

$$\eta = \frac{P_{out}}{P_{in}} \cdot 100\% = \frac{24652}{29640} \cdot 100\% = 83.2\%.$$

Example. Suppose fan blades are attached to the shaft of the above motor. This adds $15Nm$ of additional torque loss, independent of speed. What is the new speed?

The total developed torque is now

$$T_{dev} = 200 + 12 + 15 = 227Nm,$$

where $200Nm$ is from the load, and $27Nm$ is from rotational losses. The armature current increases to

$$I_A = \frac{T_{dev}}{K\phi} = \frac{227}{2.387} = 95.1$$

By KVL, we have

$$E_A = V_T - I_A R_A = 300 - (95.1)(0.065) = 293.82V.$$

Thus,

$$\omega_m = \frac{E_A}{K\phi} = \frac{293.82}{2.387} = 123.09rad/s,$$

$$n_m = \omega_m \cdot \frac{60}{2\pi} = 1175.4rpm.$$

33.2 Separately Excited DC Machines

This configuration is similar to shunt-connected, except the field and armature have separate sources. That is, the field circuit consists of a voltage source V_F from $-$ to $+$, through a resistor R_F and inductor L_F with a current I_F leading back to the voltage source. The armature circuit consists of V_T from $-$ to $+$ with current I_A through resistor R_A reaching a motor consuming E_A before leading back to the voltage source.

33.3 Permanent-Magnet DC Motors

This type of motor is similar to separately excited, except the field is produced by permanent magnets. It is useful in fractional-horsepower applications, such as for small fans, power windows, windshield wipers, and servos.

33.4 Series-Connected DC Motors

This type of motor consists of the field and armature connected in series. Thus, instead of splitting into two paths, we encounter V_T from $-$ to $+$, followed by L_F , R_F , R_A , and E_A from $+$ to $-$. Here, L_F and R_F constitute the field, while R_A and E_A form the armature. In this case, $I_A = I_F$. Series-connected motors have high torque at low speeds. They are suitable for application such as electric automotive starter motors, electrical drills, screwdrivers, and handheld mixers.

33.5 Torque-Speed Characteristics

All motors are characterized by torque-speed characteristics.

34 Monday April 10, 2017

34.1 Formulas for Final Exam

A separate formula sheet is not allowed for this course. The following machine equations and conversions will be given:

$$T_{dev} = K\phi I_A,$$

$$E_A = K\phi\omega_m,$$

$$P = T\omega_m,$$

$$1HP = 746W,$$

$$\omega_m(\text{rad/s}) = n_m(\text{rev/min}) \cdot 2\pi(\text{rads/rev}) \cdot \frac{1}{60}(\text{min/s}).$$

Some things that will not be given but should be remembered are shown below. For an inductor,

$$V(t) = L \frac{di(t)}{dt},$$

$$Z_L = j\omega L.$$

For a capacitor,

$$i(t) = C \frac{dV(t)}{dt},$$

$$Z_C = \frac{1}{j\omega C}.$$

For DC and AC power, we follow the passive reference convention. When current flows from $-$ to $+$ across a circuit element, then $P = -Vi$. When current flows from $+$ to $-$ across a circuit element, the $P = Vi$. Thus, $P > 0$ when it is absorbed, and $P < 0$ when it is delivered.

For AC power, P is the power or average power, while Q is the reactive power. They can be expressed as

$$P = \frac{V_{rms}^2}{R} = I_{rms}^2 R,$$

$$Q = \frac{V_{rms}^2}{X} = I_{rms}^2 X,$$

where $X = \|Z_L\| = \omega L$ for an inductor, and $X = \|Z_C\| = 1/\omega C$ for a capacitor. Alternatively, they are given as

$$P = I_{rms} V_{rms} \cos(\theta),$$

$$Q = I_{rms} V_{rms} \sin(\theta),$$

where $\theta = \theta_V - \theta_I$ is the power angle. Complex power is given by

$$\bar{S} = P + jQ = \frac{1}{2}\overline{VI}.$$

The power triangle where P is adjacent, Q is opposite, and P_{app} is the hypotenuse with an angle of θ between P and P_{app} can be expressed as

$$P_{app} = \sqrt{P^2 + Q^2} = \|\bar{S}\|.$$